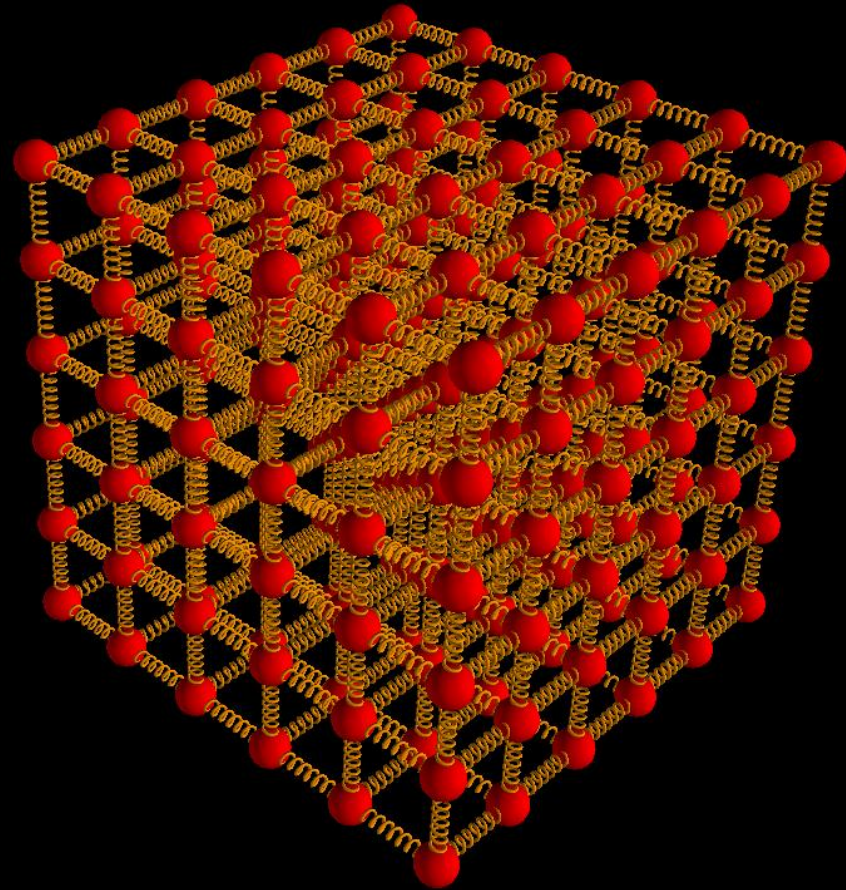
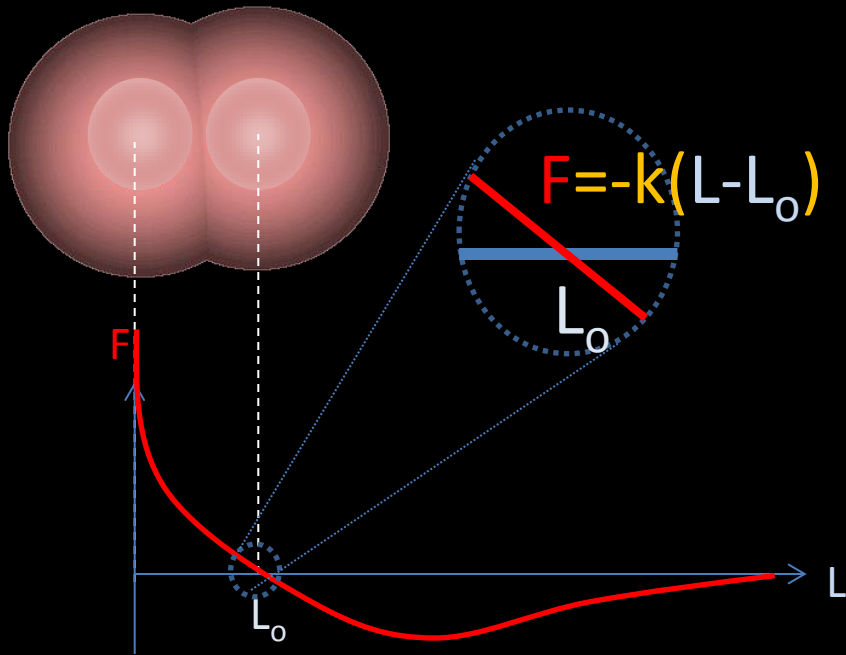


Wed.	4.6-.7, .9-.10 Stress, Strain, Young's Modulus, Compression, Sound <i>InStove @ noon      Science Poster Session: Hedco7pm~9pm</i>	RE 4.b
Lab	L4: Young's Modulus & Speed of Sound (Read 4.11-.12)	
Fri	4.11-.12; .14-.15 Sound in Solids, Analytical Solutions Quiz 3	RE 4.c
Mon.	4.8, .13 Friction and Buoyancy & Suction	RE 4.d
Tues.		EP 4, HW4: Ch 4 Pr's 46, 50, 81, 88 & CP
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Fri.	Exam 1 (Ch 1-4)	

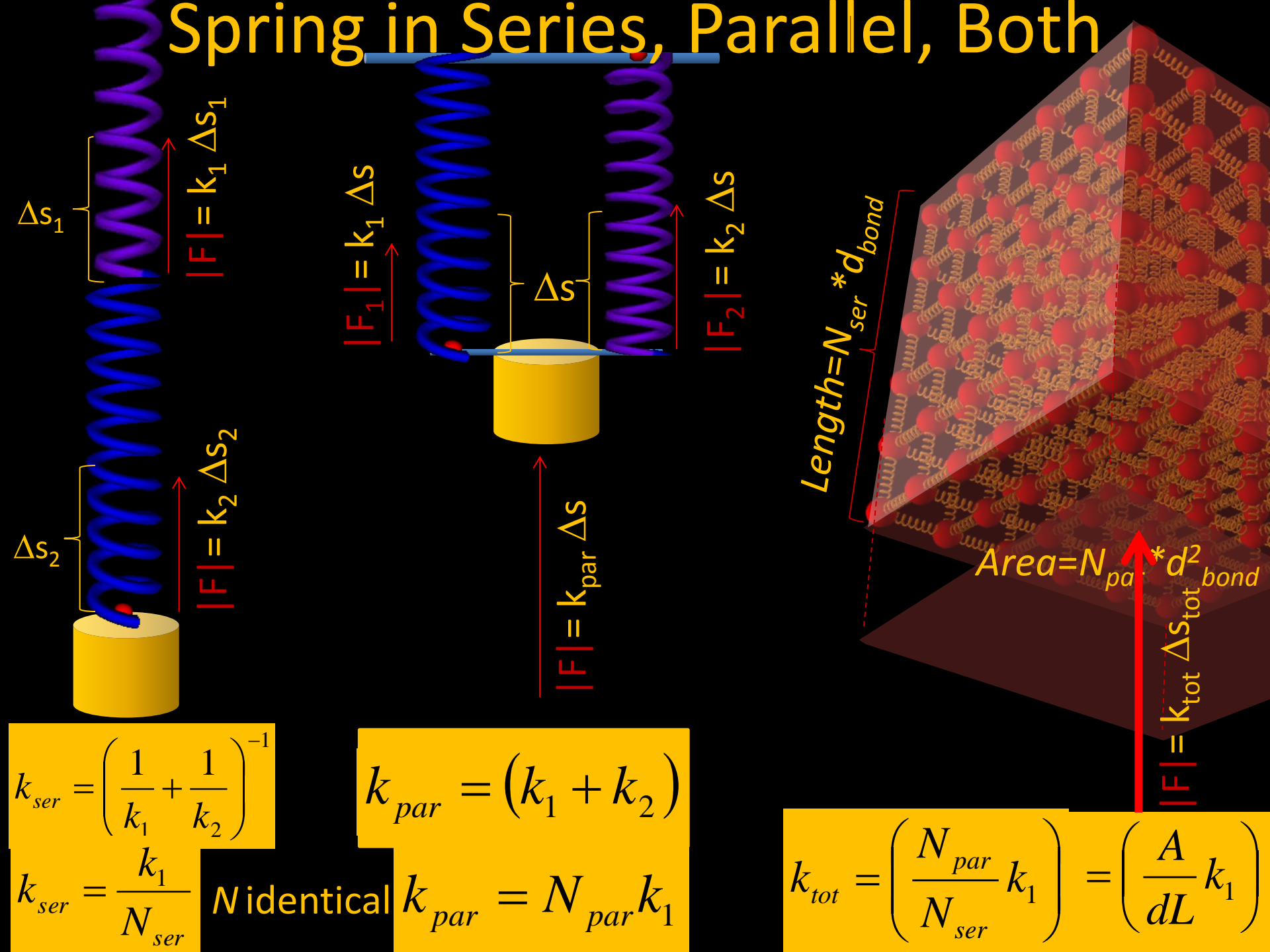
# Ball-Spring Model

Molecule

Solid



# Spring in Series, Parallel, Both



$$k_{ser} = \left( \frac{1}{k_1} + \frac{1}{k_2} \right)^{-1}$$

$$k_{ser} = \frac{k_1}{N_{ser}}$$

$N$  identical

$$k_{par} = (k_1 + k_2)$$

$$k_{par} = N_{par} k_1$$

$$k_{tot} = \left( \frac{N_{par}}{N_{ser}} k_1 \right) = \left( \frac{A}{dL} k_1 \right)$$

# Spring in Series & Parallel Rephrased Stress, Strain, and Young's Modulus

$$F = k_{total} \Delta L$$

regroup

$$F = \left( k_{atomic} \left( \frac{A}{dL} \right) \right) \Delta L$$

Microscopic details

$$F = \left( \frac{k_{atomic}}{d} \right) \frac{A \Delta L}{L} \text{ Macroscopic measurables}$$

From counting series and parallel

$Y \equiv$  Young's Modulus

$$F = Y \frac{A \Delta L}{L}$$

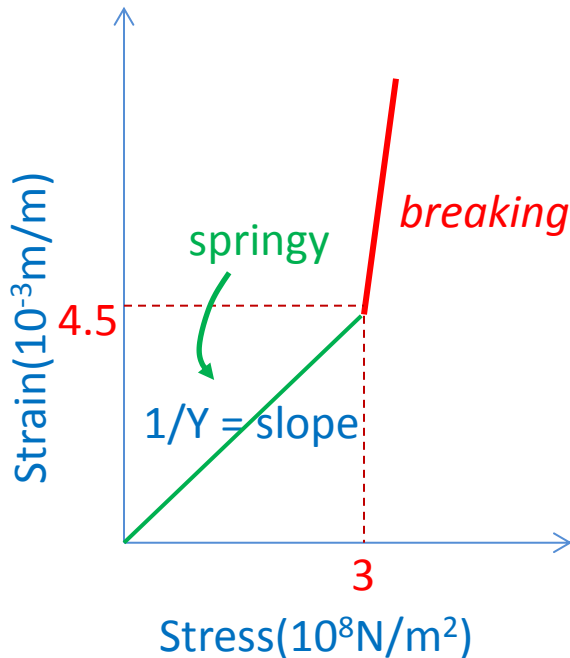
$$\left( \frac{F}{A} \right) = Y \left( \frac{\Delta L}{L} \right)$$

**Stress**

force transmitted per area (pressure)

**Strain**

stretch per length (fractional stretch)



Two wires with equal lengths are made of pure copper. The diameter of wire A is twice the diameter of wire B.

When 6 kg masses are hung on the wires, wire B stretches more than wire A.

$$Y = (F/A)/(DL/L) = k/d$$

You make careful measurements and compute Young's modulus for both wires. What do you find?

1)  $Y_A > Y_B$

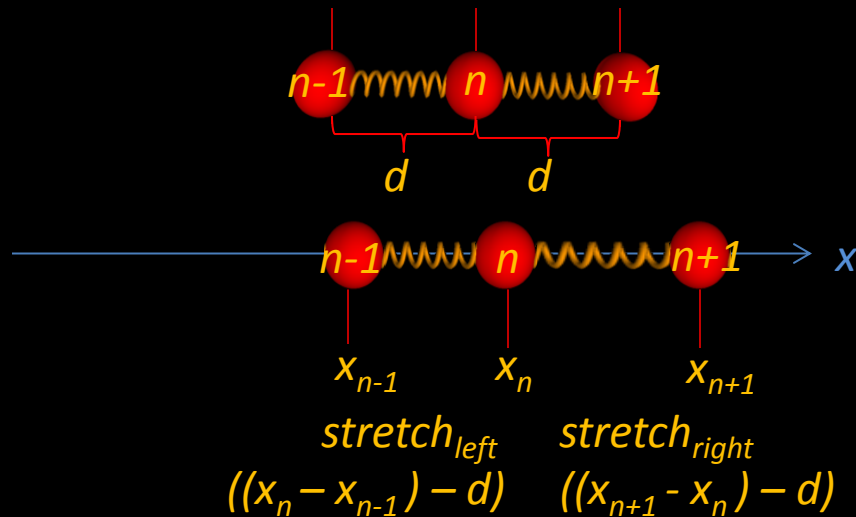
2)  $Y_A = Y_B$

3)  $Y_A < Y_B$

**Example:** You hang a heavy ball with a mass of 14 kg from a silver rod 2.6m long by 1.5 mm by 3.1mm. You measure a stretch of the rod, and find that the rod stretched 0.002898 m. Using these experimental data, what value of Young's modulus do you get?

The density of silver is  $10.5 \text{ g/cm}^3$  and you can look up its atomic mass. What's the inter-atomic spring stiffness?

# Speed of Sound in a Solid: the logic



$$F_{n,Left} = -k_s ((x_n - x_{n-1}) - d)$$

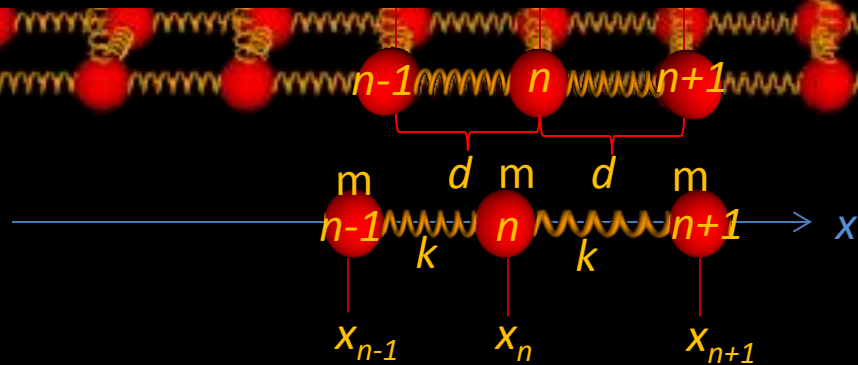
$$F_{n,Right} = k_s ((x_{n+1} - x_n) - d)$$

$$F_{n,net} = F_{n,Left} + F_{n,Right} = -k_s ((x_n - x_{n-1}) - d) + k_s ((x_{n+1} - x_n) - d)$$

$$F_{n,net} = k_s (x_{n-1} + x_{n+1} - 2x_n)$$

Copmutational simulation

# Speed of Sound in a Solid



$$F_{n,net} = k_s (x_{n-1} + x_{n+1} - 2x_n - x_{n-1})$$

## Informed Guess at $v$ 's dependence

$$v = \sqrt{\frac{k_s}{m}} d$$

Stiffer, for a given atomic displacement, greater force pulling it so greater velocity achieved.

More distance between atoms means further the distortion can propagate just through the light weight spring /bond without encountering the resistance of massive atoms.

More massive, more inertial resistance to applied force, less velocity achieved.

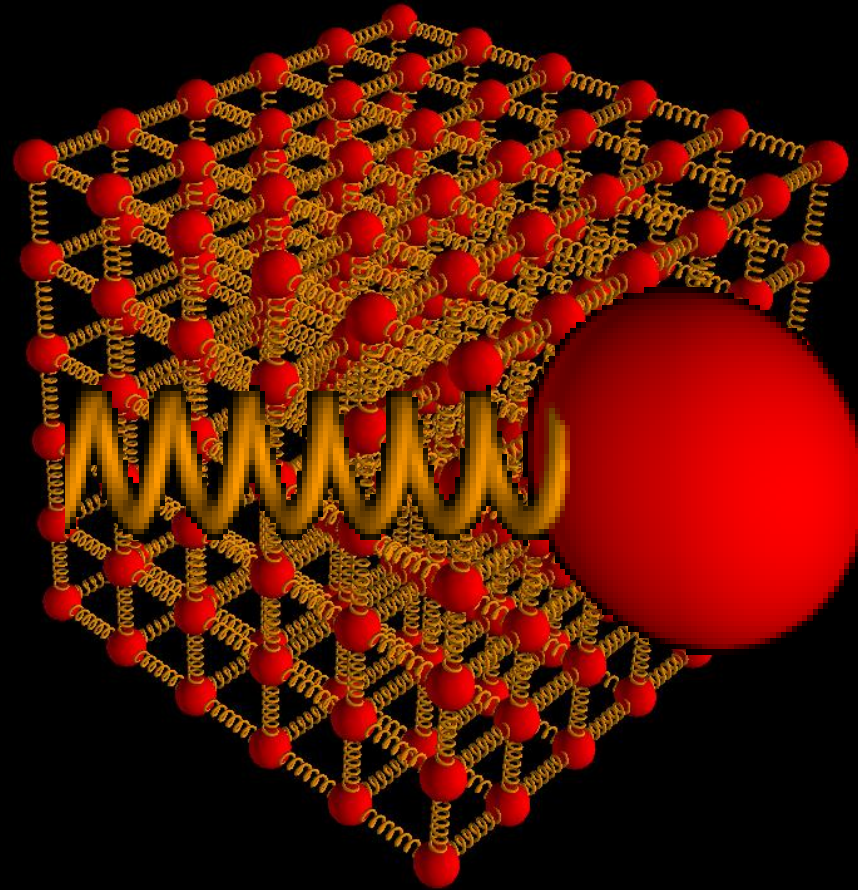
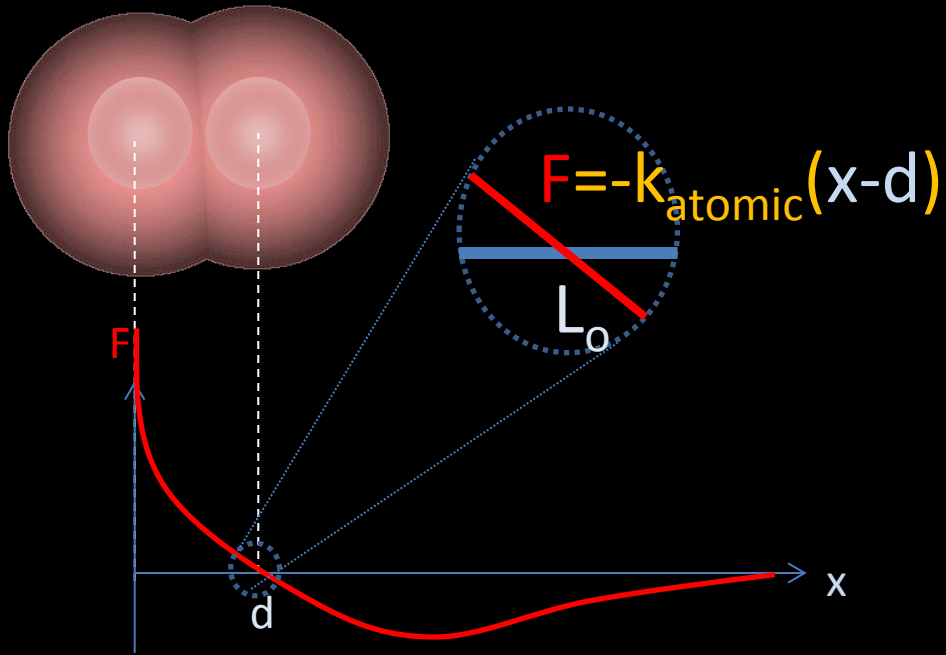


**Example:** The spring constant of aluminum is about 16 N/m. The typical separation of Al atoms was  $2.6 \times 10^{-10}$  m. Recall also that the atomic mass of aluminum is 27 g/mole. So what is the speed of sound in Aluminum?

# Microscopic to Macroscopic Springs

Molecule

Solid



$$\vec{F} = -k_{\text{spring}} \left( \begin{array}{c} \hat{\Delta \vec{L}}_i \\ \vdots \end{array} \right) \Delta \vec{L}$$

# Case Study in Three Modes of Exploration with Varying Force: Mass on Spring Experimentation / Observation

# Case Study in Three Modes of Exploration with Varying Force: Mass on Spring

## Experimentation / Observation

## Computation / Simulation

### Observations to Understand

- force, velocity, and position vary sinusoidally
- force and position vary in synch
- velocity varies out-of-synch
- Period's dependence
  - Mass - greater mass, slower
  - Stiffness - greater stiffness, faster
  - Amplitude - no effect !

### Observations to Understand

- Changing gravity only changes center of oscillation

$$\vec{F}_s = -k_s (|L| - |L_o|) \hat{L}$$

$$L\_mag = \text{mag}(\text{ball.pos})$$

$$L\_hat = \text{ball.pos} / L\_mag$$

$$F = -k * (L\_mag - L_o) * L\_hat$$

$$\vec{p}_f = \vec{p}_i + \vec{F}_s \Delta t$$

$$\text{ball.p} = \text{ball.p} + F * \text{deltat}$$

$$\vec{r}_f = \vec{r}_i + \frac{\vec{p}}{m} \Delta t$$

$$\text{ball.pos} = \text{ball.pos} + (\text{ball.p} / \text{ball.m}) * \text{deltat}$$

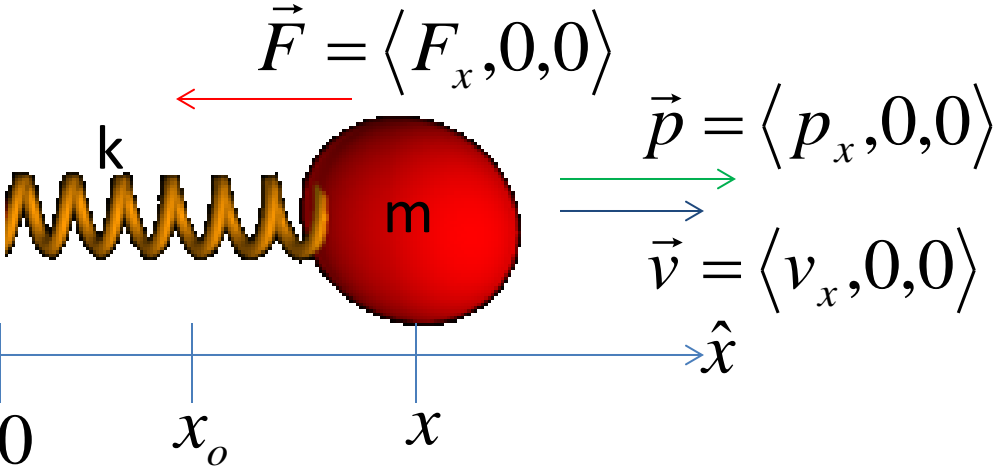
# Finite changes to infinitesimal changes: derivatives

# Case Study in Three Modes of Exploration with

## Varying Force: Mass on Spring

### Theory / Analysis

**System: Ball**



**Guess from experiment and simulation**

$$x(t) = X \cos\left(2\pi \frac{t}{T}\right) + x_o$$

**Shorthand:**  $\omega \equiv \frac{2\pi}{T}$

$$x(t) = X \cos(\omega t) + x_o$$

$$F_x(t) = -k * [x(t) - x_o]$$

$$\frac{dp_x(t)}{dt} = -k * [x(t) - x_o]$$

$$\frac{d[mv_x(t)]}{dt} = -k * [x(t) - x_o]$$

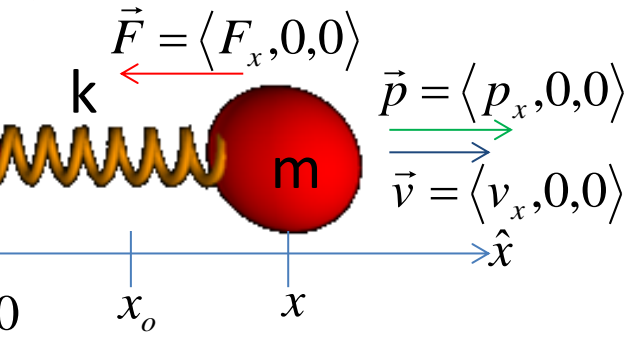
$$m \frac{d}{dt} \left[ \frac{dx(t)}{dt} \right] = -k * [x(t) - x_o]$$

$$\frac{d^2 x(t)}{dt^2} = -\frac{k}{m} * [x(t) - x_o]$$

# Case Study in Three Modes of Exploration with Varying Force: Mass on Spring

## Theory / Analysis

**System: Ball**



$$F_x(t) = -k * [x(t) - x_o]$$

$$\frac{d^2 x(t)}{dt^2} = -\frac{k}{m} * [x(t) - x_o]$$

**Guess**

$$x(t) = X \cos(\omega t) + x_o$$

where:  $\omega \equiv \frac{2\pi}{T}$

Plug in and see if guessed solution works

$$\frac{d^2}{dt^2} [X \cos(\omega t) + x_o] = -\frac{k}{m} * [X \cos(\omega t) + x_o - x_o]$$

$$\cancel{X} \frac{d^2}{dt^2} [\cancel{\cos(\omega t)}] = -\frac{k}{m} * [\cancel{X} \cancel{\cos(\omega t)}]$$

$$\frac{d}{dt} [-\omega \sin(\omega t)] = -\frac{k}{m} \cos(\omega t)$$

$$\cancel{\omega^2} \cancel{\cos(\omega t)} = \cancel{-\frac{k}{m}} \cancel{\cos(\omega t)}$$

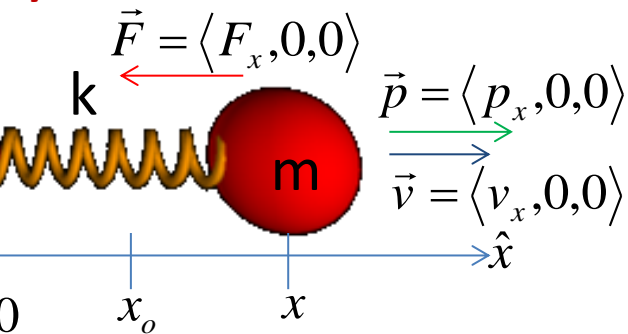
$$\omega^2 = \frac{k}{m}$$

Our guess works if  $\omega = \sqrt{\frac{k}{m}}$

# Case Study in Three Modes of Exploration with Varying Force: Mass on Spring

## Theory / Analysis

**System: Ball**



$$F_x(t) = -k * [x(t) - x_o]$$

$$\frac{d^2 x(t)}{dt^2} = -\frac{k}{m} * [x(t) - x_o]$$

**Solution**

$$x(t) = X \cos(\omega t) + x_o$$

where:  $\omega \equiv \frac{2\pi}{T} = \sqrt{\frac{k}{m}}$

**Concisely tells us...**

$$x(t) = X \cos(\omega t) + x_o$$

- Sinusoidally oscillates
- About the equilibrium
- With a period that...
  - Shortens with greater stiffness
  - Lengthens with larger masses

$$\omega \equiv \frac{2\pi}{T} = \sqrt{\frac{k}{m}} \Rightarrow T = 2\pi \sqrt{\frac{m}{k}}$$

- Doesn't care about amplitude



## Period dependence on: mass

Suppose the period of a spring-mass oscillator is 1 s. What will be the period if we double the mass?

- a.  $T = 0.5 \text{ s}$
- b.  $T = 0.7 \text{ s}$
- c.  $T = 1.0 \text{ s}$
- d.  $T = 1.4 \text{ s}$
- e.  $T = 2.0 \text{ s}$

## Period dependence on Stiffness:

Suppose the period of a spring-mass oscillator is 1 s. What will be the period if we double the spring stiffness? (We could use a stiffer spring, or we could attach the mass to two springs.)

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## Period Dependence on Amplitude:

Suppose the period of a spring-mass oscillator is 1 s with an amplitude of 5 cm. What will be the period if we increase the amplitude to 10 cm, so that the total distance traveled in one period is twice as large?

1)  $T = 0.5 \text{ s}$

2)  $T = 0.7 \text{ s}$

3)  $T = 1.0 \text{ s}$

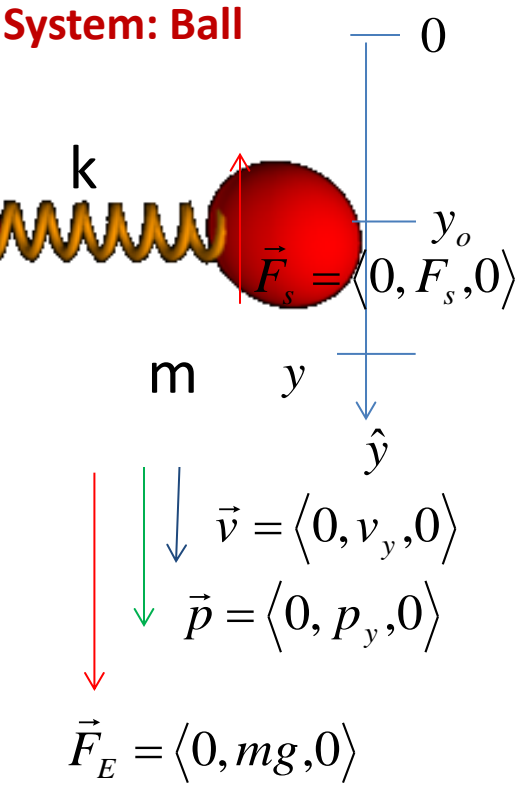
4)  $T = 1.4 \text{ s}$

5)  $T = 2.0 \text{ s}$

# Case Study in Three Modes of Exploration with Varying Force: Mass on Spring

## Theory / Analysis

How does gravitational interaction change behavior?



Note: I've defined *down* as +y direction  
So Earth's pull has + sign

$$\vec{F}_{net} = \vec{F}_s + \vec{F}_E = \langle 0, F_s + F_E, 0 \rangle$$

$$F_{net.y}(t) = -k * [y(t) - y_o] + mg$$

$$F_{net.y}(t) = -k * [y(t) - y_o] + \frac{k}{k} mg$$

$$F_{net.y}(t) = -k * [y(t) - y_o] + k \left[ \frac{mg}{k} \right]$$

$$F_{net.y}(t) = -k * \left[ y(t) - y_o - \frac{mg}{k} \right]$$

$$F_{net.y}(t) = -k * \left[ y(t) - \left\{ y_o + \frac{mg}{k} \right\} \right]$$

$$F_{net.y}(t) = -k * [y(t) - y'_o]$$

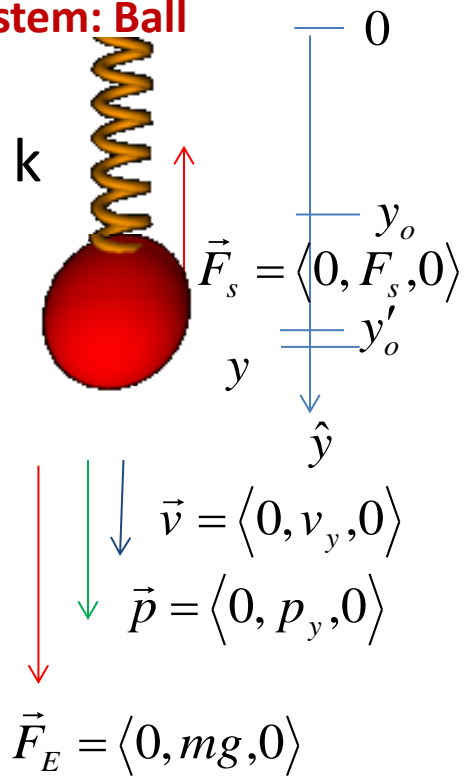
where  $y'_o \equiv y_o + \frac{mg}{k}$

# Case Study in Three Modes of Exploration with Varying Force: Mass on Spring

## Theory / Analysis

How does gravitational interaction change behavior?

System: Ball



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$$F_{net.y}(t) = -k * [y(t) - y_o] + mg$$

⋮

$$F_{net.y}(t) = -k * [y(t) - y'_o] \text{ where } y'_o \equiv y_o + \frac{mg}{k}$$

- Exact same form as for horizontal mass-spring, but shifted equilibrium

$$m \frac{d^2}{dt^2} y(t) = -k * [y(t) - y'_o]$$

Solution:

$$y(t) = Y \cos(\omega t) + y'_o$$

$$\omega \equiv \frac{2\pi}{T} = \sqrt{\frac{k}{m}} \Rightarrow T = 2\pi \sqrt{\frac{m}{k}}$$

## Period dependence on $g$ :

Suppose the period of a spring-mass oscillator is 1 s with an amplitude of 5 cm. What will be the period if we take the oscillator to a massive planet where  $g = 19.6 \text{ N/kg}$ ?

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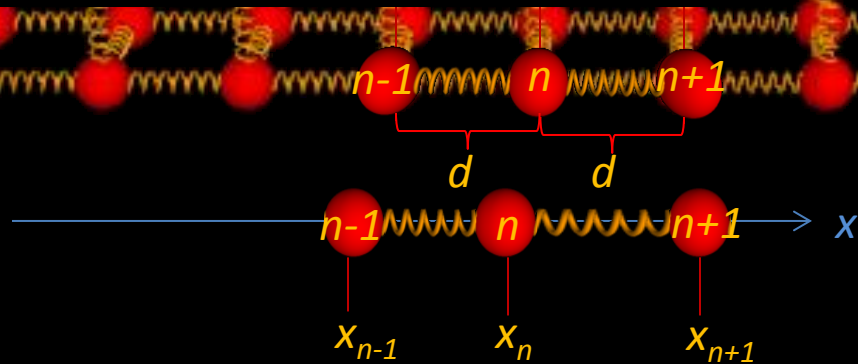
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# Speed of Sound in a Solid: the logic



$$F_{n,net} = k_s (x_{n-1} + x_{n+1} - 2x_n)$$

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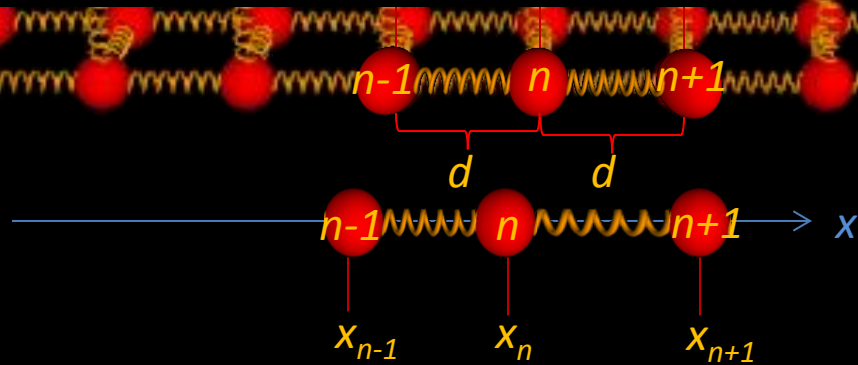
$$\frac{d\left(\frac{dx_n}{dt}\right)}{dt} = \frac{k_s}{m} (x_{n-1} + x_{n+1} - 2x_n)$$

$$\frac{d^2 x_n}{dt^2} = \frac{k_s}{m} (x_{n-1} + x_{n+1} - 2x_n)$$

**Speed of Sound in a Solid:  
the result**

$$v = \sqrt{\frac{k_s}{m}} d$$

# Speed of Sound in a Solid



Stiffer, for a given atomic displacement, greater force pulling it so greater velocity achieved.

$$v = \sqrt{\frac{k_s}{m}} d$$

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