

Wed.	(C14) 4.3 Electric Displacement Washington 3-2 Rep 7pm AHoN 116	HW5
Thurs.		
Fri.	(C14) 4.4.1 Linear Dielectrics (read rest at your discretion)	

From last Time: Polarization

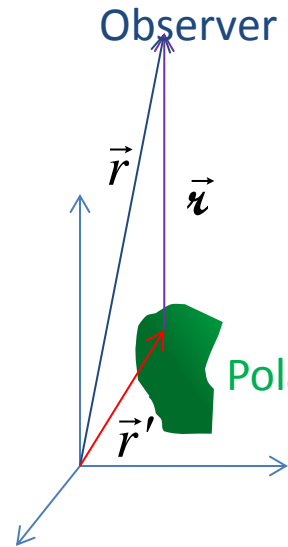
$$\vec{P} \equiv \frac{d\vec{p}}{d\tau}$$

$$V_{dips}(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int_{\text{volume}} \left(\frac{\hat{z}}{r^2} \right) \cdot \vec{P}(\vec{r}') d\tau' = \frac{1}{4\pi\epsilon_0} \left[\int_{\text{surface}} \frac{\sigma_b}{r} da' + \int_{\text{volume}} \frac{\rho_b}{r} d\tau' \right]$$

where

$$\sigma_b = \vec{P} \cdot \hat{a} \quad \text{and} \quad \rho_b = -\vec{\nabla} \cdot \vec{P}$$

$$Q_b = \int_{\text{volume}} \rho_b d\tau' + \oint_{\text{surface}} \sigma_b da$$



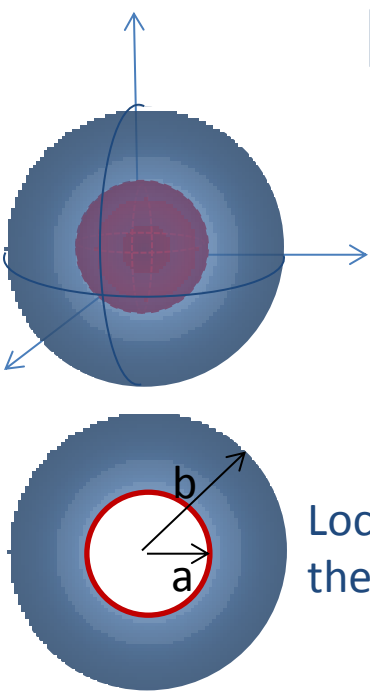
Polarization & “Bound Charge”

Exercise

$$\sigma_b = \vec{P} \cdot \hat{a} \quad \text{and} \quad \rho_b = -\vec{\nabla} \cdot \vec{P}$$

A thick spherical shell (inner radius a and outer radius b) is made of dielectric material with a “frozen-in” polarization $\vec{P}(\vec{r}) = \frac{k}{r} \hat{r}$

Locate all of the bound charge and use Gauss’s law to calculate the electric field in the three regions.



Cross-sectional view

The Electric “Displacement”

Quite Generally, Gauss’s law says

$$\vec{\nabla} \cdot \vec{E} = \frac{1}{\epsilon_0} \rho_{(all)}$$

Now we also relate “bound” charge (due to variation in density of dipoles)

$$-\vec{\nabla} \cdot \vec{P} = \rho_b$$

So, if you have a region with dipoles *and* free charges,

$$\rho_{(all)} = \rho_{free} + \rho_{bound}$$

or,

$$\rho_{free} = \rho_{(all)} - \rho_{bound}$$

$$\rho_{free} = \epsilon_0 \vec{\nabla} \cdot \vec{E} - (-\vec{\nabla} \cdot \vec{P})$$

$$\rho_{free} = \vec{\nabla} \cdot (\epsilon_0 \vec{E} + \vec{P})$$

$$\underbrace{\epsilon_0 \vec{E} + \vec{P}}_{\equiv \vec{D}} \text{ Electric Displacement}$$

So Gauss’s Law for *free* charge and Electric *Displacement*

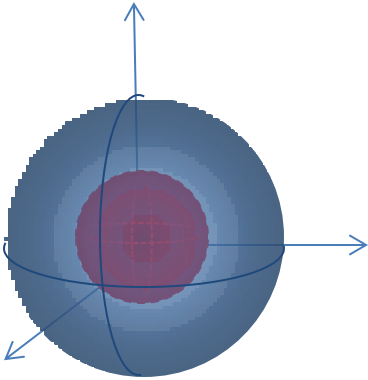
$$\rho_{free} = \vec{\nabla} \cdot \vec{D}$$

and

$$Q_{free} = \int \rho_{free} d\tau = \int \vec{D} \cdot d\vec{a}$$

Polarization & Electric Displacement

Exercise – take 2



$$Q_{free} = \int \rho_{free} d\tau = \int \vec{D} \cdot d\vec{a} \quad \epsilon_0 \vec{E} + \vec{P} \equiv \vec{D}$$

A thick spherical shell (inner radius a and outer radius b) is made of dielectric material with a “frozen-in” polarization $\vec{P}(\vec{r}) = \frac{k}{r} \hat{r}$. There are no free charges.

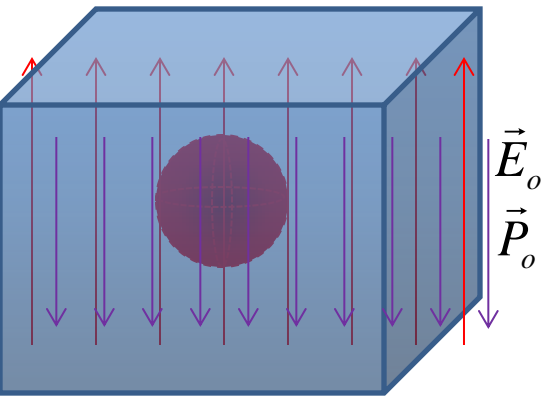
Find D from our new Gauss’s Law in all three regions, and then find E from it.

Polarization & Electric Displacement

Example

$$Q_{free} = \int \rho_{free} d\tau = \int \vec{D} \cdot d\vec{a} \quad \epsilon_0 \vec{E} + \vec{P} \equiv \vec{D}$$

Consider a huge slab of dielectric material initially with uniform field, \vec{E}_o and corresponding uniform polarization and electric displacement $\vec{D}_o = \epsilon_0 \vec{E}_o + \vec{P}_o$



You cut a small spherical hole out of it. What is the field in its center in terms of \vec{E}_o and \vec{P}_o ?

For illustrative purposes only, take the polarization to be anti-parallel to the field, and imagine both to be in the z direction.

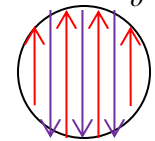
By Superposition Principle, cutting out a sphere is the same as inserting a sphere of opposite polarization.

Quoting Example 4.2 (which in turn builds on 3.9), the field *inside* a uniformly polarized sphere is

$$\vec{E}_{sphere} = -\frac{1}{3\epsilon_0} \vec{P}_{sphere}$$

So, we 'add in' a sphere of polarization $-\vec{P}_o$

Adding field $\vec{E}_{added} = \frac{1}{3\epsilon_0} \vec{P}_o$



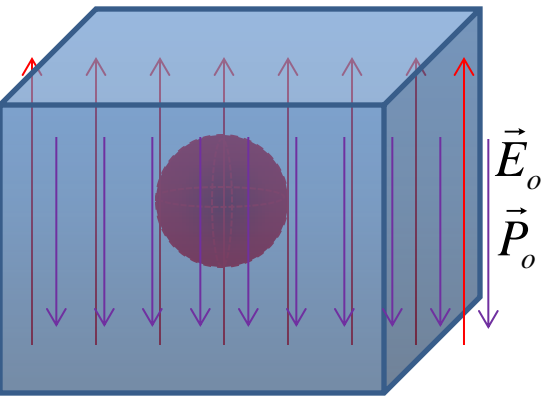
$$\vec{E}_{in.sphere} = \vec{E}_o + \frac{1}{3\epsilon_0} \vec{P}_o$$

Polarization & Electric Displacement

Example

$$Q_{free} = \int \rho_{free} d\tau = \int \vec{D} \cdot d\vec{a} \quad \epsilon_0 \vec{E} + \vec{P} \equiv \vec{D}$$

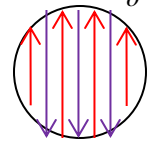
Consider a huge slab of dielectric material initially with uniform field, \vec{E}_o and corresponding uniform polarization and electric displacement $\vec{D}_o = \epsilon_0 \vec{E}_o + \vec{P}_o$



You cut a small spherical hole out of it. What is the field in its center in terms of \vec{E}_o and \vec{P}_o ?

$$\vec{E}_{in.sphere} = \vec{E}_o + \frac{1}{3\epsilon_0} \vec{P}_o$$

$$\vec{E}_{sphere} = -\frac{1}{3\epsilon_0} \vec{P}_{sphere}$$



What is the electric displacement in its center in terms of \vec{D}_o and \vec{P}_o ?

$$\vec{D}_{sphere} = \epsilon_0 \vec{E}_{sphere} + \vec{P}_{sphere}$$

There is no material in the sphere, so

$$\vec{D}_{sphere} = \epsilon_0 \left(\vec{E}_o + \frac{1}{3\epsilon_0} \vec{P}_o \right) \quad \text{where} \quad \vec{E}_o = \frac{1}{\epsilon_0} (\vec{D}_o - \vec{P}_o)$$

so

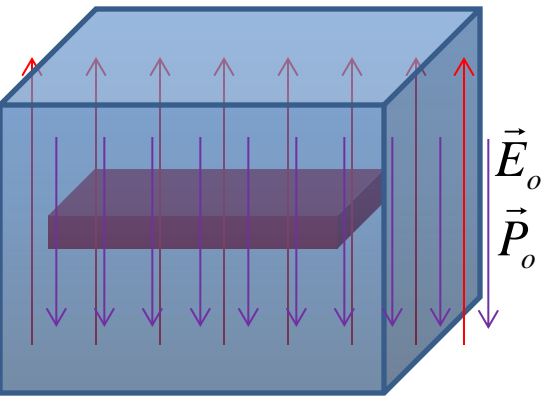
$$\vec{D}_{sphere} = \epsilon_0 \left(\frac{1}{\epsilon_0} (\vec{D}_o - \vec{P}_o) + \frac{1}{3\epsilon_0} \vec{P}_o \right) = \left(\vec{D}_o - \frac{2}{3} \vec{P}_o \right)$$

Polarization & Electric Displacement

Exercise

$$Q_{free} = \int \rho_{free} d\tau = \int \vec{D} \cdot d\vec{a} \quad \epsilon_0 \vec{E} + \vec{P} \equiv \vec{D} \quad \sigma_b = \vec{P} \cdot \hat{a} \quad \text{and} \quad \rho_b = -\vec{\nabla} \cdot \vec{P}$$

Consider a huge slab of dielectric material initially with uniform field, \vec{E}_o and corresponding uniform polarization and electric displacement $\vec{D}_o = \epsilon_0 \vec{E}_o + \vec{P}_o$.



You cut out a wafer-shaped cavity perpendicular to \vec{P}_o .

What is the field in its center in terms of \vec{E}_o and \vec{P}_o ?

Hint: Think of *inserting* the appropriate wafer-sized capacitor.

What is the electric displacement in its center in terms of \vec{D}_o and \vec{P}_o ?

The Electric “Displacement”

So Gauss’s Law for *free* charge and Electric *Displacement*

$$\rho_{free} = \vec{\nabla} \cdot \vec{D} \quad \text{and} \quad Q_{free} = \int \rho_{free} d\tau = \int \vec{D} \cdot d\vec{a} \quad \text{where} \quad \vec{D} \equiv \epsilon_0 \vec{E} + \vec{P}$$

Of practical use – with polarizable materials, you might directly control ρ_{free} , but ρ_{bound} unavoidably changes in response (reminiscent of the free energies in thermo)

Warning: while $\vec{\nabla} \times \vec{E} = 0$ in electrostatics, and so $\vec{E} = \frac{1}{4\pi\epsilon_0} \int \frac{\rho}{r^2} \hat{r} d\tau'$

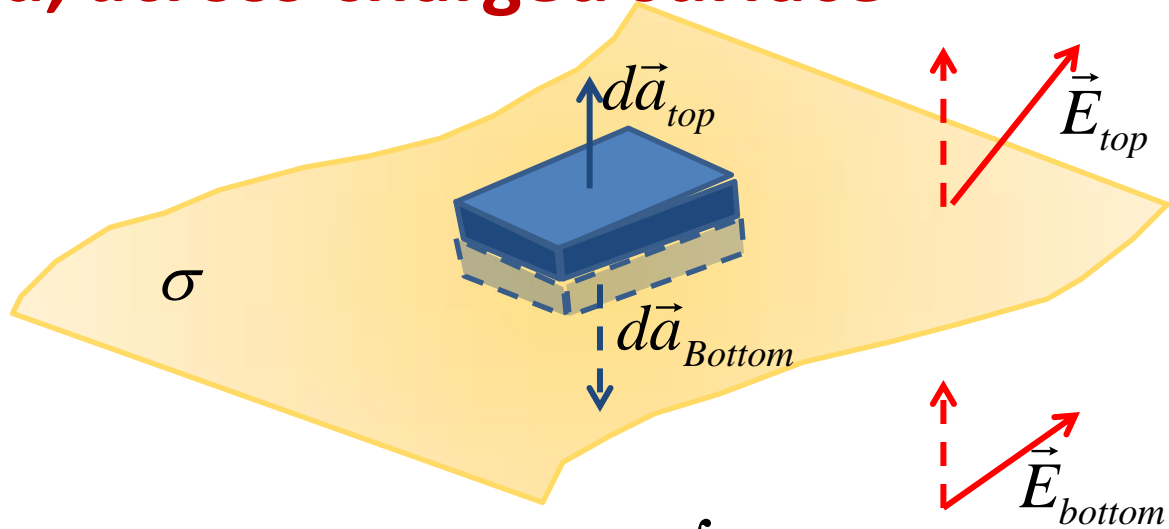
For electric displacement $\vec{\nabla} \times \vec{D} = \vec{\nabla} \times \epsilon_0 \vec{E} + \underbrace{\vec{\nabla} \times \vec{P}}$

$$\text{and so } \vec{D} \neq \frac{1}{4\pi\epsilon_0} \int \frac{\rho_{free}}{r^2} \hat{r} d\tau'$$

Not necessarily 0
Which means it can’t necessarily be expressed as gradient of a scalar field

Boundary Conditions

Electric field, *across* charged surface



$$\vec{\nabla} \cdot \vec{E} = \frac{\rho_{encl}}{\epsilon_0}$$

$$\oint \vec{E} \cdot d\vec{a} = \frac{Q_{encl}}{\epsilon_0}$$

$$\int \vec{E}_{top} \cdot d\vec{a}_{top} + \int \vec{E}_{bottom} \cdot d\vec{a}_{bottom} + \int \vec{E}_{sides} \cdot d\vec{a}_{sides} = \frac{Q_{encl}}{\epsilon_0} = \frac{\int \sigma da_{surface}}{\epsilon_0}$$

Send side height / area to 0

$$\int \vec{E}_{top} \cdot d\vec{a}_{top} + \int \vec{E}_{bottom} \cdot d\vec{a}_{bottom} = \frac{\int \sigma da_{surface}}{\epsilon_0}$$

$$E_{\perp top} A + E_{\perp bottom} A(-1) = \frac{\sigma A}{\epsilon_0}$$

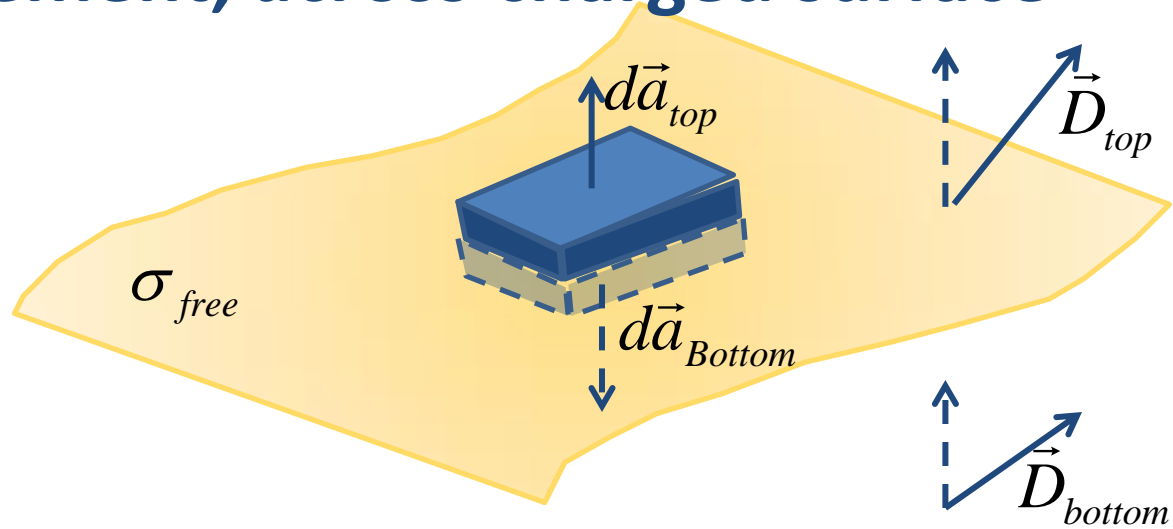
$$E_{\perp top} - E_{\perp bottom} = \frac{\sigma}{\epsilon_0}$$

Boundary Conditions

Electric Displacement, *across* charged surface

$$\vec{\nabla} \cdot \vec{D} = \rho_{free.encl}$$

$$\oint \vec{D} \cdot d\vec{a} = Q_{free.encl}$$



$$\int \vec{D}_{top} \cdot d\vec{a}_{top} + \int \vec{D}_{bottom} \cdot d\vec{a}_{bottom} + \int \vec{D}_{sides} \cdot d\vec{a}_{sides} = Q_{free.encl} = \int \sigma_{free} da_{surface}$$

Send side height / area to 0

$$\int \vec{D}_{top} \cdot d\vec{a}_{top} + \int \vec{D}_{bottom} \cdot d\vec{a}_{bottom} = \int \sigma_{free} da_{surface}$$

$$D_{\perp top} A + D_{\perp bottom} A(-1) = \sigma_{free} A$$

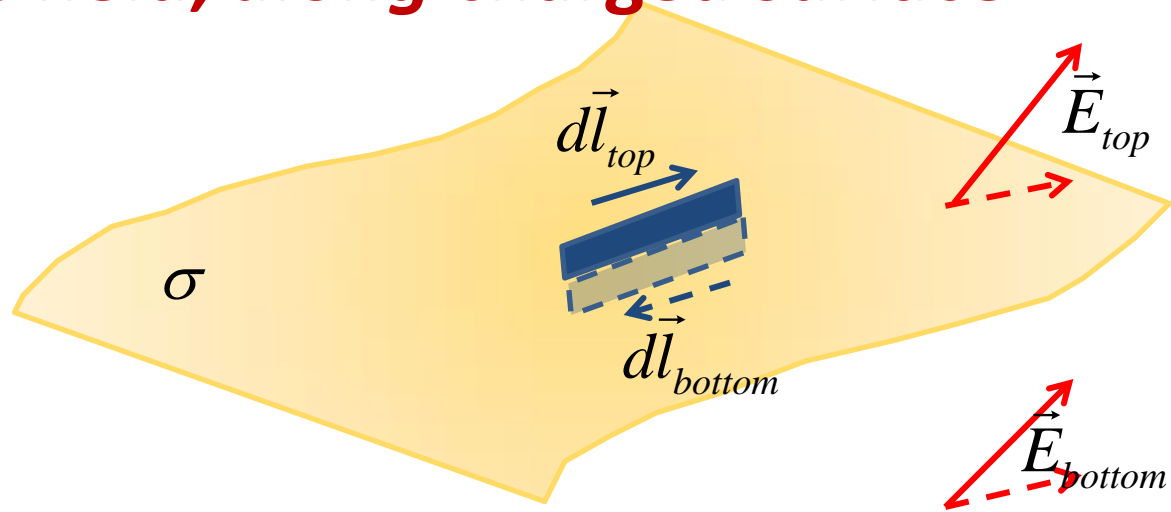
$$D_{\perp top} - D_{\perp bottom} = \sigma_{free}$$

Boundary Conditions (static) Electric field, *along* charged surface

$$\vec{\nabla} \times \vec{E} = 0$$

$$\oint \vec{E} \cdot d\vec{l} = 0$$

$$\int \vec{E}_{top} \cdot d\vec{l}_{top} + \int \vec{E}_{bottom} \cdot d\vec{l}_{bottom} + \int \vec{E}_{sides} \cdot d\vec{l}_{sides} = 0$$



Send side height to 0

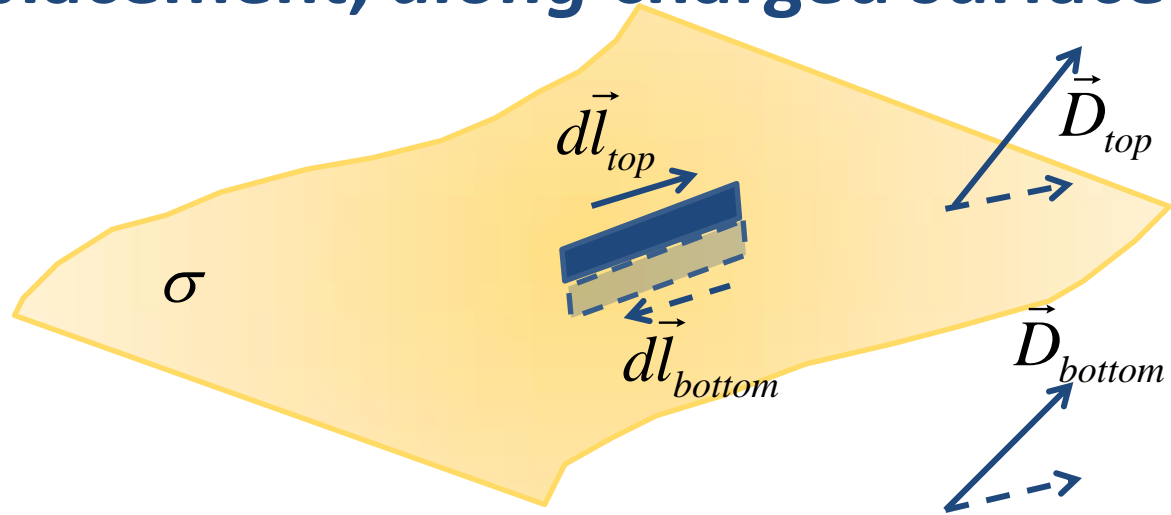
$$\int \vec{E}_{top} \cdot d\vec{l}_{top} + \int \vec{E}_{bottom} \cdot d\vec{l}_{bottom} = 0$$

$$E_{\parallel top} L + E_{\parallel bottom} L(-1) = 0$$

$$E_{\parallel top} - E_{\parallel bottom} = 0$$

Boundary Conditions

(static) Electric displacement, *along* charged surface



$$\vec{\nabla} \times \vec{D} = \vec{\nabla} \times \vec{P}$$

$$\oint \vec{D} \cdot d\vec{l} = \oint \vec{P} \cdot d\vec{l}$$

$$\int \vec{D}_{top} \cdot d\vec{l}_{top} + \int \vec{D}_{bottom} \cdot d\vec{l}_{bottom} + \int \vec{D}_{sides} \cdot d\vec{l}_{sides} = \int \vec{P}_{top} \cdot d\vec{l}_{top} + \int \vec{P}_{bottom} \cdot d\vec{l}_{bottom} + \int \vec{P}_{sides} \cdot d\vec{l}_{sides}$$

Send side height to 0

$$\int \vec{D}_{top} \cdot d\vec{l}_{top} + \int \vec{D}_{bottom} \cdot d\vec{l}_{bottom} = \int \vec{P}_{top} \cdot d\vec{l}_{top} + \int \vec{P}_{bottom} \cdot d\vec{l}_{bottom}$$

$$D_{\parallel top} L + D_{\parallel bottom} L(-1) = P_{\parallel top} L + P_{\parallel bottom} L(-1)$$

$$D_{\parallel top} - D_{\parallel bottom} = P_{\parallel top} - P_{\parallel bottom}$$

Boundary Conditions Electric and Displacement fields

Along

$$E_{\parallel top} - E_{\parallel bottom} = 0$$

$$D_{\parallel top} - D_{\parallel bottom} = P_{\parallel top} - P_{\parallel bottom}$$

(could have guessed as much from $\vec{D} = \epsilon_0 \vec{E} + \vec{P}$.)

Across

$$E_{\perp top} - E_{\perp bottom} = \frac{\sigma}{\epsilon_0}$$

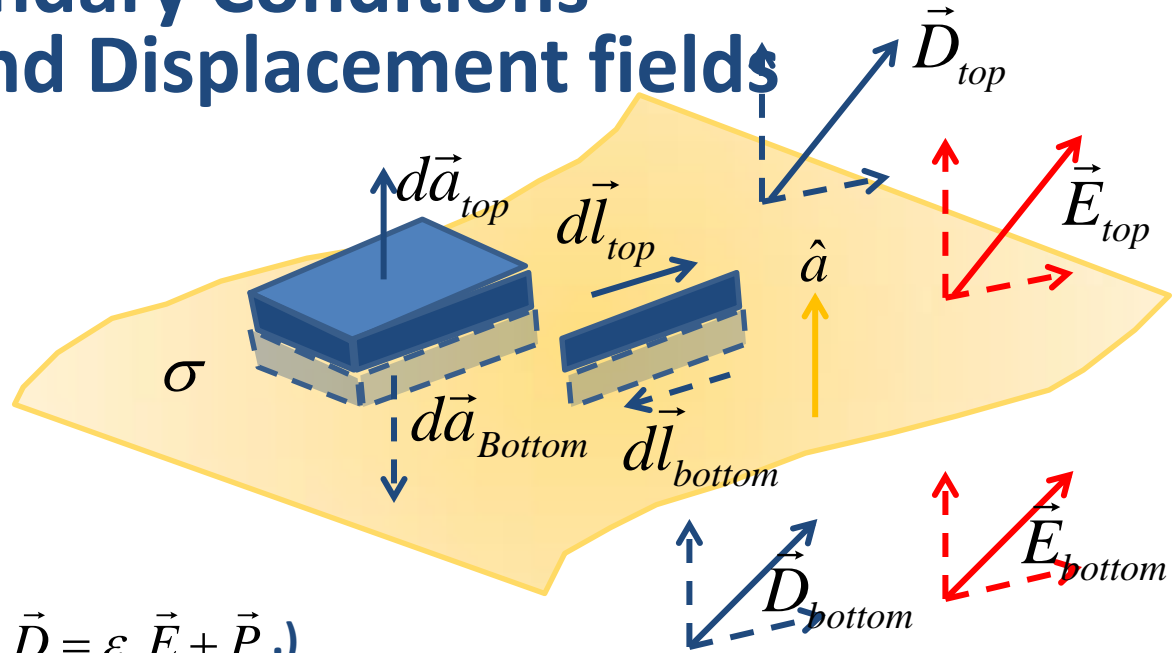
or

$$D_{\perp top} - D_{\perp bottom} = \sigma_{free}$$

$$\vec{E}_{top} \cdot \hat{a} - \vec{E}_{bottom} \cdot \hat{a} = \frac{\sigma}{\epsilon_0}$$

$$\vec{D}_{top} \cdot \hat{a} - \vec{D}_{bottom} \cdot \hat{a} = \sigma_{free}$$

(could have guessed as much from $\vec{D} = \epsilon_0 \vec{E} + \vec{P}$ and $\sigma - \sigma_{free} = \sigma_b = \vec{P} \cdot \hat{a}$.)



Boundary Conditions Electric and Displacement fields

Along

$$E_{\parallel top} - E_{\parallel bottom} = 0$$

$$D_{\parallel top} - D_{\parallel bottom} = P_{\parallel top} - P_{\parallel bottom}$$

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P}$$

$$\sigma - \sigma_{free} = \sigma_b = \vec{P} \cdot \hat{a}$$

Across

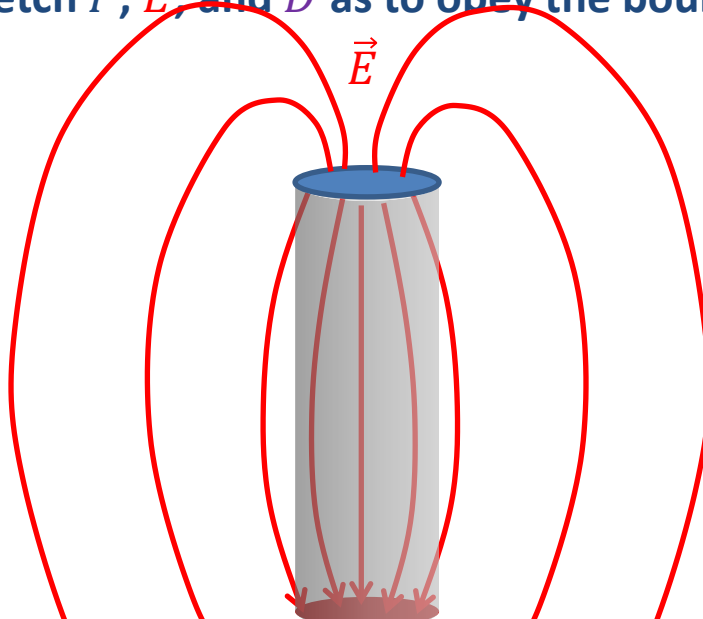
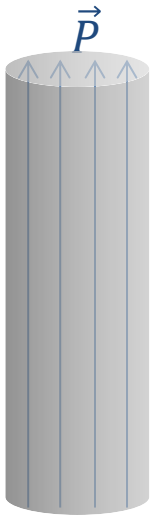
$$E_{\perp top} - E_{\perp bottom} = \frac{\sigma}{\epsilon_0}$$

$$D_{\perp top} - D_{\perp bottom} = \sigma_{free}$$

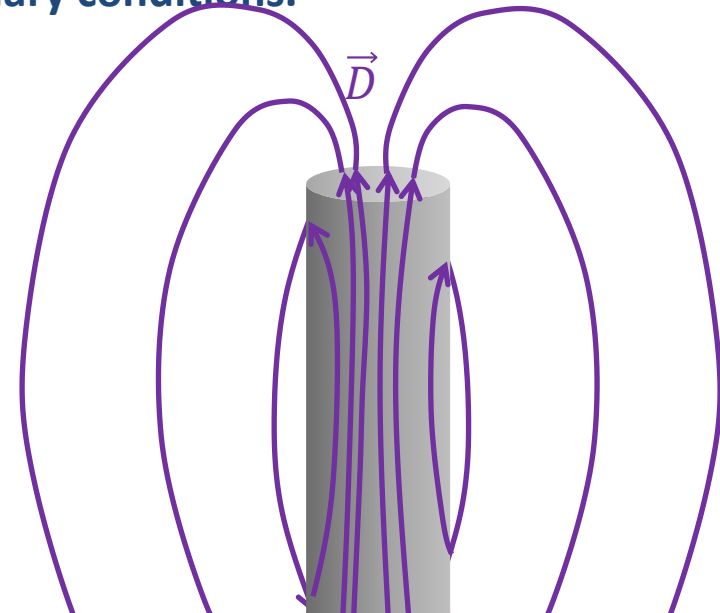
Exercise

Bar Electret (like an electric bar magnet): uniform \vec{P} along axis

Sketch \vec{P} , \vec{E} , and \vec{D} as to obey the boundary conditions.



Only surface charge is the bound surface charge on two faces



There is no free surface charge, so no discontinuity in D . $\vec{D} = \epsilon_0 \vec{E} + \vec{P}$

Mon.	(C14) 4.3 Electric Displacement	
Mon.	(C14) 4.4.1 Linear Dielectrics (read rest at your discretion)	
Wed.	6.1 Magnetization	HW10