Fri.	(C 14) 4.1 Polarization	
Mon. Wed.	(C 14) 4.2 Field of Polarized Object (C14) 4.3 Electric Displacement Washington 3-2 Rep. 7pm AHoN 116	HW5
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For *small* stretch, first term in Taylor Series (Hook's Law)



Atom's Response to Electric Field $\vec{p} \approx \alpha E_{ant}$ How Small is Small Enough? Special case: uniform 'electron cloud' $\rho = \frac{q}{\frac{4}{2}\pi a^3}$ (easier to consider nucleus being displaced) Atom on a stick $\vec{F}_{ext \to nuc} = q\vec{E}_{ext \to nuc} = -q\vec{E}_{elect \to nuc}$ $\vec{E}_{ext \to nuc} = -\vec{E}_{elect \to nuc}$ $\vec{E}_{external}$ a $\oint \vec{E}_{external} = q \vec{E}_{external} \qquad \oint \vec{E}_{elect \to nuc} \cdot d\vec{a} = \frac{1}{\varepsilon} \int \rho d\tau' = \frac{1}{\varepsilon_{o}} \int \frac{-q}{\frac{4}{3}\pi a^{3}} d\tau$ $\vec{F}_{internal} = E_{elect \to nuc} 4\pi d^{2} = \frac{-q}{\frac{4}{2}\pi a^{3}\varepsilon} \frac{4}{3}\pi d^{3}$ $E_{elect \to nuc} = \frac{-q}{4\pi a^3 \varepsilon d^2} d^3 = \frac{-q}{4\pi a^3 \varepsilon} d$ $-\vec{E}_{ext\to nuc} = \vec{E}_{elect\to nuc} = \frac{-q}{4\pi a^3 \varepsilon_o} d\vec{d}$ $4\pi a^3 \varepsilon_o \vec{E}_{ext\to nuc} = q d\hat{d}$ Key to this working: charge distribution is radially uniform $\alpha \vec{E}_{avt \rightarrow nuc} = \vec{p}$ and angularly symmetric

Atom's Response to Electric Field Key to this working: charge $\vec{p} \approx \alpha E_{axt}$ distribution is radially uniform How Small is Small Enough?and angularly symmetric Real (simple): electron cloud $\rho(r) = \frac{q}{\pi a^3} e^{-2r/a}$ (easier to consider nucleus being displaced) $\vec{F}_{ovt \rightarrow nuc} = q\vec{E}_{ovt \rightarrow nuc} = -q\vec{E}_{elect \rightarrow nuc}$ Atom on a stick $E_{ext \rightarrow nuc} = -E_{elect \rightarrow nuc}$ **E**_{external} $\oint \vec{E}_{elect \to nuc} \cdot d\vec{a} = \frac{1}{\varepsilon} \int \rho d\tau' = \frac{1}{\varepsilon_{o}} \int \frac{-q}{\pi a^{3}} e^{-2r/a} d\tau$ $e_{rnal} = q E_{external}$ $E_{elect \to nuc} 4\pi d^2 = \frac{1}{\varepsilon_{e}} \frac{-q}{\pi a^3} 4\pi \int_{0}^{u} e^{-2r/a} r^2 dr$ $\dot{F}_{internal} = -$ **F**_{external} Assuming d<<a $E_{elect \to nuc} 4\pi d^2 \approx \frac{1}{\varepsilon} \frac{-q}{\pi a^3} 4\pi \int_{0}^{a} (1 - \frac{2r}{a}) r^2 dr \quad e^{-2r/a} \approx 1 - \frac{2r}{a}$ Indeed, constant polarizability only for radially constant charge $E_{elect \to nuc} 4\pi d^2 \approx \frac{1}{\varepsilon} \frac{-q}{\pi a^3} 4\pi \left(\frac{1}{3} d^3 - \frac{1}{2} \frac{d^4}{a}\right)$ distribution $E_{elect \to nuc} \approx \frac{1}{\varepsilon} \frac{-q}{\pi a^3} d\left(\frac{1}{3} - \frac{1}{2}\frac{d}{a}\right)$ But if d<<a, good enough approximation $\varepsilon_o \pi a^3 E_{elect \to nuc} \approx -qd\left(\frac{1}{3} - \frac{1}{2}\frac{d}{a}\right) = -p\left(\frac{1}{2} - \frac{1}{2}\frac{d}{a}\right)$

Atom's Response to Electric Field $\vec{p} \approx \vec{\alpha} \vec{E}_{ext}$

Polarizability Tensor

Mass on un-even springs analogy

Atom's Response to Electric Field $\vec{p} \approx \vec{\alpha} \vec{E}_{ext}$

Polarizability Tensor

 $\vec{E} = E\hat{z}$

 $= \alpha_{r}E_{r}$ Imagine cylindrical 'molecule' which is easier to polarize along length than radially: $\alpha_{l} > \alpha_{r}$. So when field is applied in z-direction, Get poloarization at an angle

$$\begin{bmatrix} p_r \\ p_l \end{bmatrix} \approx \begin{bmatrix} \alpha_r & 0 \\ 0 & \alpha_l \end{bmatrix} \begin{bmatrix} E_r \\ E_l \end{bmatrix}$$

More generally expressed in terms of Cartesian coordinates that may not be aligned with p or E,

$$p_{x} = \alpha_{xx}E_{x} + \alpha_{xy}E_{y} + \alpha_{xz}E_{z}$$

$$p_{y} = \alpha_{yx}E_{x} + \alpha_{yy}E_{y} + \alpha_{yz}E_{z}$$

$$p_{z} = \alpha_{zx}E_{x} + \alpha_{zy}E_{y} + \alpha_{zz}E_{z}$$

$$\vec{p} = \begin{bmatrix} \alpha_{xx} & \alpha_{xy} & \alpha_{xz} \\ \alpha_{yx} & \alpha_{yy} & \alpha_{yz} \\ \alpha_{zx} & \alpha_{zy} & \alpha_{zz} \end{bmatrix} \begin{bmatrix} E_{x} \\ E_{y} \\ E_{z} \end{bmatrix}$$

$$\vec{p} = \vec{\alpha}\vec{E}$$



Response of a dipole to fields

Torque in uniform field



 $N = pE\sin\theta$

 \vec{E}_{ext}

 $\vec{p} = q\vec{s}$

s/2

Change in energy with Rotation

 S_i

$$\Delta U(\vec{r}) = -\int_{\vec{r}_i}^{\vec{r}_f} \vec{F} \cdot d\vec{\ell}$$

$$\Delta U_{rot} = -2\int_{\vec{r}_i}^{\vec{r}_f} q\vec{E} \cdot \frac{s}{2} d\vec{\theta} = -\int_{\vec{r}_i}^{\vec{r}_f} qEs \sin\theta d\theta = -qsE(\cos\theta_f - \cos\theta_i)$$

$$\Delta U_{rot} = -\Delta(\vec{p} \cdot \vec{E})|_{p,E}$$

Change in energy with vibration

$$\Delta U_{stretch} = -2\int_{\vec{r}_i}^{s_f} q\vec{E} \cdot \frac{d\vec{s}}{2} = -2\int_{\vec{r}_i}^{s_f} q\vec{E} \cdot d\vec{s} = -\Delta(\vec{p} \cdot \vec{E})|_{E,\theta}$$

 S_i

Response of a dipole to fields $\vec{N} = \vec{p} \times \vec{E}$ $N = pE \sin \theta$

Exercise: A dipole is a distance *z* above an infinite grounded conducting plane. The dipole makes an angle θ with the perpendicular to the plane. Find the torque on it. In what orientation is it stable?



Response of a dipole to fields

Force in uniform field

 $\vec{F}_{net} = \sum \vec{F}_i$

р

g

+C

 $= q\vec{s}$

$$\vec{F}_{net} = q\vec{E} - q\vec{E} = 0$$

Force in a non-uniform field

$$\begin{split} \vec{F}_{net} &= q\vec{E}_{+} - q\vec{E}_{-} \\ \vec{F}_{net} &= q\Big(\vec{E}\big(\vec{r}_{+}^{\prime}\big) - \vec{E}\big(\vec{r}_{-}^{\prime}\big)\Big) \\ \vec{F}_{net} &= q\Big(\vec{E}\big(\vec{r}_{+}^{\prime}\big) - \vec{E}\big(\vec{r}_{+}^{\prime} - \vec{s}\,\big)\Big) \end{split}$$

First, imagine dipole along z-axis,

$$\vec{F}_{net} = qs \left(\frac{\vec{E}(\vec{r}_{+}') - \vec{E}(\vec{r}_{+}' - \vec{s})}{s} \right) = p_z \left(\frac{\partial \vec{E}}{\partial z} \right)$$
Alternatively, if dipole points $p_y \left(\frac{\partial \vec{E}}{\partial y} \right)$ or dipole points $p_x \left(\frac{\partial \vec{E}}{\partial x} \right)$
generally

$$\vec{F}_{net} = p_x \left(\frac{\partial \vec{E}}{\partial x}\right) + p_y \left(\frac{\partial \vec{E}}{\partial y}\right) + p_z \left(\frac{\partial \vec{E}}{\partial z}\right) = \left(\vec{p} \cdot \vec{\nabla}\right) \vec{E}$$



Response of an *induced* dipole to fields

 $\vec{F}_{net} = \left(\vec{p} \cdot \vec{\nabla} \right) \vec{E}$ and $\vec{p} = \vec{\alpha} \vec{E}$

$$\vec{F}_{net} = \left(\vec{\alpha}\vec{E}\cdot\vec{\nabla}\right)\vec{E}$$

SO

Product rule 4 (if polarizability is uniform / scalar)

$$\vec{F}_{net} = \alpha \left(\frac{1}{2} \vec{\nabla} \left(E^2 \right) - \vec{E} \times \left(\vec{\nabla} \times \vec{E} \right) \right)$$

Only if time-varying current and/or charge distribution

Note: force in direction of increasing field strength - "Optical tweezers"

Bulk material's Response to Electric Field Polarization

$\mathbf{P} \equiv$ dipole moment per unit volume



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