| Mon. | (C 17) 12.1.1-.1.2, 12.3.1 E to B; 5.1.1-.1.2 Lorentz Force Law: fields and forces |  |
| :--- | :--- | :--- |
| Wed | (C 17) 5.1.3 Lorentz Force Law: currents |  |
| Thurs. |  |  |
| Fri. | (C 17) 5.2 Biot-Savart Law | HW6 |

## Force between stationary charges

(Coulomb's Law: Eq’n 2.1)


$$
\vec{F}_{q \rightarrow Q}=\frac{1}{4 \pi \varepsilon_{o}} \frac{q Q}{r_{q \rightarrow Q}^{2}} \hat{r}=\frac{q Q}{4 \pi \varepsilon_{o}} \frac{\vec{r}}{(r)^{3}}
$$

## Force between moving charges



$$
\vec{u} \equiv c \hat{r}-\vec{v}
$$

$$
\vec{F}_{Q \leftarrow q}=\frac{q Q}{4 \pi \varepsilon_{o}} \frac{\tau}{(\vec{\imath} \cdot \vec{u})^{3}}\left\{\left[\left(c^{2}-v^{2}\right) \vec{u}+\vec{\imath} \times(\vec{u} \times \vec{a})\right]+\frac{\vec{V}}{c} \times\left[\hat{\imath} \times\left[\left(c^{2}-v^{2}\right) \vec{u}+\vec{\imath} \times(\vec{u} \times \vec{a})\right]\right\}\right.
$$

"The entire theory of classical electrodynamics is contained in that equation...but you see why I preferred to start out with Coulomb's law." - Griffiths

$$
\begin{aligned}
& \vec{F}_{Q \leftarrow q}=\frac{q Q}{4 \pi \varepsilon_{o}} \frac{r}{(\vec{r} \cdot \vec{u})^{3}}\{\underbrace{\left[\left(c^{2}-v^{2}\right) \vec{u}+\vec{r} \times(\vec{u} \times \vec{a})\right]}_{\text {Electric }}+\frac{\vec{v}}{\frac{V}{c} \times\left[\hat{r} \times\left[\left(c^{2}-v^{2}\right) \vec{u}+\vec{r} \times(\vec{u} \times \vec{a})\right]\right]}\} \\
& \text { Depends on observer's } \\
& \text { perception of source } \\
& \text { charge's velocity and } \\
& \text { acceleration } \\
& \text { Also depends on } \\
& \text { observer's perception of } \\
& \text { recipient charge's velocity }
\end{aligned}
$$

But who's moving how fast is all relative

Note: extremely asymmetric between two charges; reciprocity (Newton's $3^{\text {rd }}$ ) does not hold

## Crash Course in Special Relativity

Principle: Laws of Physics should be the same in all inertial frames of reference

laws of E\&M referenced an absolute speed, not speed relative to preferred frame
Observations:
Charge moving near stationary magnet feels magnetic force; charge stationary near moving magnet feels...?

One of the three must go

Key to new frame-transformation math:
Speed of light is measured to be the
 same in all reference frames

## Crash Course in Special Relativity

New frame-transformation math derived - light clocks and meter sticks
with clock Measured in frame $\Delta t_{w . c}=2 \frac{h}{c} \backslash d=\sqrt{(2 h)^{2}+L_{W . S}^{2}}=\sqrt{(2 h)^{2}+\left(v \Delta t_{w . S}\right)^{2}}$ at rest with two events defining interval


$$
L_{w . c}=v \Delta t_{w . c}
$$

With meter stick
Measured in frame at

$$
L_{w . c}=v \Delta t_{w . S} \sqrt{1-\left(\frac{v}{c}\right)^{2}}
$$

rest with two locations of defining events.

$$
\begin{aligned}
\frac{d}{\Delta t_{W . S}} \Delta t_{W . S} & =\sqrt{\left(c \Delta t_{w . c}\right)^{2}+\left(v \Delta t_{W . S}\right)^{2}} \\
c \Delta t_{W . S} & =\sqrt{\left(c \Delta t_{w . c}\right)^{2}+\left(v \Delta t_{w . S}\right)^{2}}
\end{aligned}
$$

$$
\left(c \Delta t_{w . S}\right)^{2}=\left(c \Delta t_{w . c}\right)^{2}+\left(v \Delta t_{w . S}\right)^{2}
$$

$$
L_{w . c}=L_{w, s} \sqrt{1-\left(\frac{v}{c}\right)^{2}}
$$

$$
\left(1-\left(\frac{v}{c}\right)^{2}\right)\left(\Delta t_{w . S}\right)^{2}=\left(\Delta t_{w . c}\right)^{2}
$$

$$
L_{w, s}=\frac{L_{w, c}}{\sqrt{1-\left(\frac{v}{c}\right)^{2}}}=\gamma L_{w, c}
$$

Agreeing about both c and v means disagreeing about t and L .

## Crash Course in Special Relativity

New frame-transformation math derived - light clocks and meter sticks
Example

$$
L_{w, s}=\frac{L_{w, c}}{\sqrt{1-\left(\frac{v}{c}\right)^{2}}}=\gamma L_{w, c}
$$

$$
\Delta t_{w . c}=\Delta t_{w, S} \sqrt{1-\left(\frac{v}{c}\right)^{2}}=\frac{1}{\gamma} \Delta t_{w . S}
$$

Think of the ladder and barn problem. Say the Farmer's got hold of the ladder and he's going to run at the barn.

# Crash Course in Special Relativity Galilean Transformation Corrected 



I'm in a train watching a pool game. I see the ball role a distance $\Delta x_{\text {ball.table }}$ to the pocket. I see it's taking time $\Delta t_{\text {ball.table }}$ to get there. Meanwhile, the train and I are rolling through the station in the same direction at speed $v$. You, standing in the station, see all this happen. Classically, you'd imagine that with my yard stick, I'd measure the ball rolling a distance

$$
" \Delta x_{\text {ball.table }} "=\Delta x_{\text {ball.station }}-\mathrm{v} \Delta t_{\text {ball.station }}
$$

But to phrase that in terms of what $/$ in the train would measure, " $\Delta x_{\text {ball.table }} "=\frac{1}{\gamma} \Delta \mathrm{x}_{\text {ball.table }}$

$$
\text { so, } \frac{1}{\gamma} \Delta x_{\text {ball.table }}=\Delta x_{\text {ball. station }}-\mathrm{v} \Delta t_{\text {ball.station }} \mathrm{or}, \Delta x_{\text {ball. table }}=\gamma\left(\Delta x_{\text {ball.station }}-\mathrm{v} \Delta t_{\text {ball. station }}\right)
$$

Then again, I'd imagine with your yardstick you'd measure the ball rolling a distance

$$
" \Delta x_{\text {ball.station }} "=\Delta x_{\text {ball table }}+\mathbf{v} \Delta t_{\text {ball.table }}
$$

But in terms of what you in the station would measure, " $\Delta x_{\text {ball.station }}=\frac{1}{\gamma} \Delta \mathrm{x}_{\text {ball.station }}$
so, $\quad \Delta x$ ball $_{\text {station }} \gamma\left(\Delta x_{\text {ball table }}+\mathrm{v} \Delta t_{\text {ball.table }}\right) \quad$ Note: both observers measure distances as 'proper' lengths
Consistent only if $\Delta t_{\text {ball.station }}=\gamma\left(\Delta t_{\text {ball, table }}+\frac{v}{c^{2}} \Delta x_{\text {ball, table }}\right)$

$$
\Delta t_{\text {ball, table }}=\gamma\left(\Delta t_{\text {ball.station }}-\frac{v}{c^{2}} \Delta x_{\text {ball. station }}\right)
$$ of their sticks, but neither is a equidistant to the two eventsneither measures proper time

# Crash Course in Special Relativity Galilean Transformation Corrected 



I'm in a train watching a pool game. I see the ball role a distance $\Delta x_{\text {ball.table }}$ to the pocket. I see it's taking time $\Delta t_{\text {ball.table }}$ to get there. Meanwhile, the train and I are rolling through the station in the same direction at speed $\boldsymbol{v}$. You, standing in the station, see all this happen.

$$
v_{\text {ball.station }}=\frac{\Delta x_{\text {ball.station }}}{\Delta t_{\text {ball.station }}}=\frac{\gamma\left(\Delta x_{\text {ball.table }}+v \Delta t_{\text {ball.table }}\right)}{\gamma\left(\Delta t_{\text {ball.table }}+\frac{v}{c^{2}} \Delta x_{\text {ball.table }}\right)}
$$

Some algebra later,

$$
v_{\text {ball. station }}=\frac{v_{\text {ball table }}+\mathbf{v}}{1+\frac{v v_{\text {ball table }}}{c^{2}}}
$$

Or renaming $\mathbf{v}=\mathbf{v}_{\text {table.station }}$

$$
v_{\text {ball. station }}=\frac{v_{\text {ball table }}+v_{\text {table station }}}{1+\frac{v_{\text {table.station }} v_{\text {ball.table }}}{c^{2}}}
$$

## Crash Course in Special Relativity

New frame-transformation math derived - light clocks and meter sticks Example

$$
v_{\text {ball. station }}=\frac{v_{\text {ball table }}+v_{\text {table station }}}{1+\frac{v_{\text {table.station }} \mathcal{V}_{\text {ball.table }}}{c^{2}}}
$$

Lab frame: you and I see an electrically neutral wire (all be it, with the electrons moving)
$\lambda_{+}=\frac{e}{\Delta x_{l a b}}$
$\lambda_{-}=-\frac{e}{\Delta x_{l a b}}$
$\lambda_{+}=$ions ${ }^{\Delta x^{\prime}}$ charge density $\quad \lambda_{-}=-\lambda_{+}=$electrons' charge density (coulombs/meter) (coulombs/meter)

- = ionic atomic core
- = electron
$\mathrm{v}_{\mathrm{e}}=$ electron velocity measured by us in the "lab frame"
Stationary charge

$$
\begin{aligned}
& \vec{F}_{q \leftarrow \text { wire }}=q \vec{E}_{\text {wire }} \\
& \qquad \vec{E}_{\text {wire }}=\vec{E}_{+}+\vec{E}_{-}=\frac{1}{4 \pi \varepsilon_{o}} \frac{2 \lambda_{+}}{r}+\frac{1}{4 \pi \varepsilon_{o}} \frac{2 \lambda_{-}}{r}=\frac{1}{4 \pi \varepsilon_{o}} \frac{2 \lambda_{+}}{r}+\frac{1}{4 \pi \varepsilon_{o}} \frac{-2 \lambda_{+}}{r}=0
\end{aligned}
$$

Lab frame: you and I see an electrically neutral wire (all be it, with the electrons moving)
$\lambda_{+}=\frac{e}{\Delta x_{\text {lab }}} \quad \lambda_{-}=-\frac{e}{\Delta x_{l a b}}$
$\lambda_{+}=$ions ${ }_{\text {lab }}$ charge density $\quad \lambda_{-}=-\lambda_{+} \xlongequal{\Delta n_{a b}}$ electrons' charge density (coulombs/meter) (coulombs/meter)
$\mathrm{v}_{\mathrm{e}}=$ electron velocity measured by us in the "lab frame"
Moving charge

$$
\mathrm{q} \stackrel{\mathrm{v}_{\mathrm{q}}}{\longrightarrow} \quad \vec{F}_{q \leftarrow \text { wire }}=\text { ? }
$$

## Transition / Transformation from E to M

Charge's frame:
Chain of positive atoms moving backwards at $v_{\text {atoms }}=-v_{q}$
So spacing seen by charge is related to their stationary, "proper" separation (as seen in lab) by

$$
\Delta x_{\text {atom }}^{\prime}=\frac{\Delta x_{\text {atom. proper }}}{\gamma_{q^{\prime}}}=\frac{\Delta x_{\text {lab }}}{\gamma_{q^{\prime}}}=\Delta x_{\text {lab }} \sqrt{1-\left(\frac{v_{q}}{c}\right)^{2}}
$$

Or charge density appears compressed to

$$
\lambda_{+}^{\prime}=\frac{\gamma_{q^{\prime}} e}{\Delta x_{l a b}}
$$

Chain of electrons moving forward at only $v_{e}^{\prime}=\frac{v_{e}-v_{q}}{1-\frac{v_{e} v_{q}}{c^{2}}}$
So spacing seen by charge is related to their stationary, "proper" separation (not seen in lab) by

$$
\Delta x_{e}^{\prime}=\frac{\Delta x_{e . p r o p e r}}{\gamma_{e^{\prime}}} \quad \text { where } \quad \gamma_{e^{\prime}}=\frac{1}{\sqrt{1-\left(\frac{v_{e}^{\prime}}{c}\right)^{2}}}
$$

Similarly, separation in lab frame relates to "proper" separation (seen in electrons' rest frame)

$$
\begin{array}{r}
\Delta x_{e . p r o p e r}=\gamma_{e} \Delta x_{\text {lab }} \text { where } \gamma_{e}=\frac{1}{\sqrt{1-\left(\frac{v_{e}}{c}\right)^{2}}} \quad \begin{array}{r}
\text { Combined: } \Delta x_{e}^{\prime}=\frac{\gamma_{e} \Delta x_{l a b}}{\gamma_{e^{\prime}}} \\
\text { So, } \lambda_{-}^{\prime}=\frac{-\gamma_{e^{e}}}{\gamma_{e} \Delta x_{l a b}}=\frac{-\gamma_{e^{\prime}}}{\gamma_{e} \gamma_{q^{\prime}}} \lambda_{+}^{\prime}
\end{array}
\end{array}
$$

## Transition / Transformation from E to M

Charge's frame:


A bit of algebra later, $E_{\text {wire }}^{\prime}=\frac{1}{4 \pi \varepsilon_{o}} \frac{2 \lambda_{+} \gamma_{q^{\prime}}}{r} \frac{v_{e} v_{q}}{c^{2}} \quad$ Define $I=\lambda_{+} v_{e} \quad$ and $\quad \mu_{o} \equiv \frac{1}{\varepsilon_{o} c^{2}} \quad$ so $\quad E_{\text {wire }}^{\prime}=\frac{\mu_{o}}{4 \pi} \frac{2 I \gamma_{q^{\prime}}}{r} v_{q}$

$$
\vec{F}_{q \leftarrow \text { wire }}^{\prime}=q v_{q} \frac{\mu_{o}}{4 \pi} \frac{2 I}{r} \gamma_{q^{\prime}} \quad \text { Hand waving: } \mathrm{F}^{\sim} \text { distance/time }{ }^{2}
$$

Transformation from rest frame requires factor of $\gamma / \gamma^{2}=1 / \gamma$

Finally, we observe:

$$
\vec{F}_{q \leftarrow \text { wire }}=q v_{q} \frac{\mu_{o}}{4 \pi} \frac{2 I}{r}
$$

## Jumping in to Magnetism $\underset{\vec{u} \equiv c i-\vec{v}}{ }$

$$
\vec{F}_{Q \leftarrow q}=\frac{q Q}{4 \pi \varepsilon_{o}} \frac{r}{(\vec{r} \cdot \vec{u})^{3}}\{\underbrace{\left[\left(c^{2}-v^{2}\right) \vec{u}+\vec{r} \times(\vec{u} \times \vec{a})\right]}_{\text {Electric }}+\underbrace{\frac{\vec{V}}{c} \times\left[\hat{r} \times\left[\left(c^{2}-v^{2}\right) \vec{u}+\vec{r} \times(\vec{u} \times \vec{a})\right]\right]}_{\text {Magnetic }}\}
$$

Depends on observer's perception of source charge's velocity and acceleration

Also depends on observer's perception of recipient charge's velocity

Magnetic force

$$
\vec{F}_{Q \leftarrow q . \operatorname{mag}}=Q \vec{V} \times \underbrace{\frac{q}{4 \pi c \varepsilon_{o}} \frac{r}{(\vec{r} \cdot \vec{u})^{3}}\left\{\left[\hat{\imath} \times\left[\left(c^{2}-v^{2}\right) \vec{u}+\vec{r} \times(\vec{u} \times \vec{a})\right]\right]\right\}}_{\text {Magnetic Field, } \mathrm{B}}
$$

$$
\begin{aligned}
& \text { Cyclotron Motion in a Uniform Magnetic Field } \\
& \frac{d \vec{p}}{d t}=\vec{F}=q(\vec{E}+\vec{v} \times \vec{B}) \\
& \left.m\left\{\begin{array}{l}
d v_{x} / d t \\
d v_{y} / d t \\
d v_{z} / d t
\end{array}\right\}=q\left(\begin{array}{c}
E_{x} \\
E_{y} \\
E_{z}
\end{array}\right\}+\left\lvert\, \begin{array}{ccc}
\hat{x} & \hat{y} & \hat{z} \\
v_{x} & v_{y} & v_{z} \\
\boldsymbol{B}_{x} & \boldsymbol{B}_{y} & \boldsymbol{B}_{z}
\end{array}\right.\right)=q\left\{\begin{array}{l}
E_{x}+v_{y} \boldsymbol{B}_{z}-v_{z} \boldsymbol{B}_{y} \\
E_{y}+v_{z} \boldsymbol{B}_{x}-v_{x} \boldsymbol{B}_{z} \\
E_{z}+v_{x} \boldsymbol{B}_{y}-v_{y} B_{x}
\end{array}\right\} \\
& \text { Book's Example } \\
& m \frac{d \vec{v}}{d t}=q(\vec{E}+\vec{v} \times \vec{B}) \quad \text { In general } \\
& \text { Guess Solution Forms } \\
& \begin{aligned}
\frac{d v_{y}}{d t} & =\frac{q B}{m} v_{z} \\
\frac{d^{2} v_{y}}{d t^{2}} & =\frac{q B}{m} \frac{d v_{z}}{d t} \underbrace{\frac{d v_{z}}{d t}}_{\text {Cross next derivative }}=\frac{d^{2} v_{z}}{d t^{2}}=-\frac{q B}{m} \frac{q B}{m} v_{y} \\
\frac{d^{2} v_{y}}{d t^{2}} & =\frac{q B}{m}\left(\frac{q E}{m}-\frac{q B}{m} v_{y}\right)_{\text {Substitute }} \frac{d^{2} v_{z}}{d t^{2}}=-\left(\frac{q B}{m}\right)^{2} v_{z}
\end{aligned} \\
& \omega\left(-C_{3} \sin (\omega t)+C_{4} \cos (\omega t)\right)=\frac{q B}{m}\left(C_{1} \cos (\omega t)+C_{2} \sin (\omega t)\right) \\
& \frac{d^{2} v_{y}}{d t^{2}}=\frac{q^{2} B E}{m^{2}}-\left(\frac{q B}{m}\right)^{2} v_{y} \quad \text { Compare terms and conclude } \\
& \omega=\frac{q B}{m} \quad-C_{3}=C_{2} \quad C_{4}=C_{1}
\end{aligned}
$$

Cyclotron Motion in a Uniform Magnetic Field

$$
\vec{F}=q(\vec{E}+\vec{v} \times \vec{B})
$$

$$
v_{z}(t)=C_{1} \cos (\omega t)+C_{2} \sin (\omega t) \quad v_{y}(t)=C_{1} \sin (\omega t)-C_{2} \cos (\omega t)+\frac{E}{B}
$$

$$
\text { where } \omega=\frac{q B}{m}
$$

For position components, integrate
$z(t)=\frac{C_{1}}{\omega} \sin (\omega t)-\frac{C_{2}}{\omega} \cos (\omega t)+C_{5} \quad y(t)=-\frac{C_{1}}{\omega} \cos (\omega t)-\frac{C_{2}}{\omega} \sin (\omega t)+\frac{E}{B} t+C_{6}$

$$
\begin{array}{rr}
z(0)=-\frac{C_{2}}{\omega}+C_{5}=0 & \text { Start at origin } \quad y(0)=-\frac{C_{1}}{\omega}+C_{6}=0 \\
C_{5}=\frac{C_{2}}{\omega} & C_{6}=\frac{C_{1}}{\omega}
\end{array}
$$

$z(t)=\frac{C_{1}}{\omega} \sin (\omega t)+\frac{C_{2}}{\omega}(1-\cos (\omega t))$

$$
y(t)=\frac{C_{1}}{\omega}(1-\cos (\omega t))-\frac{C_{2}}{\omega} \sin (\omega t)+\frac{E}{B} t
$$

Cyclotron Motion in a Uniform Magnetic Field

$$
\vec{F}=q(\vec{E}+\vec{v} \times \vec{B})
$$

# $v_{z}(t)=C_{1} \cos (\omega t)+C_{2} \sin (\omega t) \quad v_{y}(t)=C_{1} \sin (\omega t)-C_{2} \cos (\omega t)+\frac{E}{B}$ 

 where $\omega=\frac{q B}{m}$
## Impose Initial Conditions

Initial position: $(0,0)$
$z(t)=\frac{C_{1}}{\omega} \sin (\omega t)+\frac{C_{2}}{\omega}(1-\cos (\omega t))$

$$
y(t)=\frac{C_{1}}{\omega}(1-\cos (\omega t))-\frac{C_{2}}{\omega} \sin (\omega t)+\frac{E}{B} t
$$

Initial Velocity: $\vec{v}(0)=0$
$v_{z}(0)=C_{1} \cos (0)+C_{2} \sin (0)$
$v_{z}(0)=C_{1}=0$
So.

$$
\begin{aligned}
& v_{y}(0)=C_{1} \sin (0)-C_{2} \cos (0)+\frac{E}{B} \\
& v_{y}(0)=-C_{2}+\frac{E}{B}=0 \quad C_{2}=\frac{E}{B}
\end{aligned}
$$

$$
v_{z}(t)=\frac{E}{B} \sin (\omega t)
$$

$$
z(t)=\frac{E}{\omega B}(1-\cos (\omega t))
$$

$$
\begin{aligned}
& v_{y}(t)=\frac{E}{B}(1-\cos (\omega t)) \\
& y(t)=\frac{E}{\omega B}(\omega t-\sin (\omega t))
\end{aligned}
$$

Cyclotron Motion in a Uniform Magnetic Field

$$
\vec{F}=q(\vec{E}+\vec{v} \times \vec{B})
$$

$$
v_{z}(t)=C_{1} \cos (\omega t)+C_{2} \sin (\omega t) \quad v_{y}(t)=C_{1} \sin (\omega t)-C_{2} \cos (\omega t)+\frac{m E}{\omega} \quad \vec{B}
$$

$$
\text { where } \omega=\frac{q B}{m}
$$

Impose Initial Conditions
Initial position: $(0,0)$

$z(t)=\frac{C_{1}}{\omega} \sin (\omega t)+\frac{C_{2}}{\omega}(1-\cos (\omega t))$

$$
v_{z}(0)=C_{1} \cos (0)+C_{2} \sin (0)
$$

$$
v_{z}(0)=C_{1}=0
$$

$$
\begin{aligned}
& y(t)=\frac{C_{1}}{\omega}(1-\cos (\omega t))-\frac{C_{2}}{\omega} \sin (\omega t)+\frac{m E}{\omega} t \\
& v_{y}(0)=C_{1} \sin (0)-C_{2} \cos (0)+\frac{m E}{\omega} \\
& v_{y}(0)=-C_{2}+\frac{m E}{\omega}=\frac{E}{B} \\
& C_{2}=\frac{m E}{\omega}-\frac{E}{B}
\end{aligned}
$$

| Mon. | (C 17) 12.1.1-.1.2, 12.3.1 E to B; 5.1.1-.1.2 Lorentz Force Law: fields and forces |  |
| :--- | :--- | :--- |
| Wed | (C 17) 5.1.3 Lorentz Force Law: currents |  |
| Thurs. |  |  |
| Fri. | (C 17) 5.2 Biot-Savart Law | HW6 |

Cyclotron Motion in a Uniform Magnetic Field

$$
\vec{F}_{\text {mag }}=\vec{F}_{B}=Q \vec{v} \times \vec{B} \quad \frac{d \vec{p}}{d t}=\vec{F}_{\text {mag }}
$$

Cyclotron Motion in a Uniform Magnetic Field

$$
\vec{F}_{\text {mag }}=\vec{F}_{B}=Q \vec{v} \times \vec{B}
$$

$\vec{B}$

$$
\begin{aligned}
& v_{x}(t)=-C_{2} \cos (\omega t)+C_{1} \sin (\omega t) \\
& v_{y}(t)=C_{1} \cos (\omega t)+C_{2} \sin (\omega t)
\end{aligned} \quad \omega=\frac{q B_{z}}{m}
$$

Integrate for positions:

$$
\begin{aligned}
& \int_{x_{o}}^{x} d x=\int_{o}^{t} v_{x}(t) d t \\
& x(t)-x_{o}=-\frac{C_{2}}{\omega} \sin (\omega t)-\frac{C_{1}}{\omega}(\cos (\omega t)-1)
\end{aligned}
$$

similarly:

$$
y(t)-y_{o}=\frac{C_{1}}{\omega} \sin (\omega t)-\frac{C_{2}}{\omega}(\cos (\omega t)-1)
$$

