Fri.	3.2 Images T4 Relaxation Method	
Mon. Wed. Thurs	3.4.14.2 Multipole Expansion 3.4.34.4 Multipole Expansion	HW4

Poisson's & Laplace's Equations

$$\vec{\nabla} \cdot \vec{E} = \rho / \varepsilon_0$$
$$\vec{E} = -\vec{\nabla} V$$

Poisson's $\nabla^2 V = -\rho/\mathcal{E}_0$

Laplace's $\nabla^2 V = 0$ In "free space" where no charges are

In Cartesian

$$\nabla^2 V = \vec{\nabla} \cdot \left(\vec{\nabla} V\right) = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0$$

Properties of Laplace's

•

$$\nabla^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0$$

No local min or max; either flat (possible tipped plane) or saddle points

 V at mid point of range is average of V around edges (proof forthcoming)

In 2D
$$V(x, y) = \frac{\oint V dl_{circle}}{2\pi R}$$

• In 3D
$$V(x, y, z) = \frac{\oint V da_{sphere}}{4\pi R^2}$$





Uniqueness Theorems (by hook or crook)

Voltages: (as with any differential equation) Regardless of *how* you've found it, if you've found one solution to Laplace / Poisson's equation that satisfies the boundary conditions, you've found the *only* solution.

Fields: Given a charge density in a cavity within a charged conductor, the field within the conductor is uniquely determined by the *inner* charge distribution and the conductor's charge amount.

Hooks and Crook: Interesting ways of finding V and E

Images: replace a problem with a simpler equivalent one (based on corollary of the first uniqueness theorem)

Relaxation: a computational method based on the potential at a point being the average of the values at the same distance (more about Next Time).

Multipole Expansion: a method for getting approximate answers for *V* far from a charge distribution (section 3.4)

Images: replace a problem with a simpler equivalent one (based on corollary of the first uniqueness theorem – if your solution works on the boundary, it works everywhere)

Example: a charge q_o suspended distance z_o above a flat conducting surface that's held at voltage V_s . Find expression for V anywhere above the conducting surface.



Example: a charge q_o suspended distance z_o above a flat conducting surface that's held at voltage V_s . Find expression for V anywhere above the conducting surface.

In gory detail (for the experience), determine image's charge and location

V(x,y,0) =

 \vec{r}_{o}

H

Â.

qo

 \vec{r}_{o}

 \vec{r}_i

 Z_i

q

θ

 $\eta \theta_{i}$

$$\begin{array}{l} \theta_o = \theta_i = 90^\circ & \frac{q_o}{\sqrt{r^2 + {z'}^2}} + \frac{q_i}{\sqrt{r^2 + {z'}^2}} = 4\pi\varepsilon_o V_s \end{array}$$

For observation location in plane of the conductor

Must be true for *all* location in plane, so can choose easyto-evaluate locations to determine values of q_i and z_i.

$$r \to \infty \quad \frac{q_o}{\sqrt{\infty^2 + {z'_o}^2}} + \frac{q_i}{\sqrt{\infty^2 + {z'_i}^2}} = 4\pi\varepsilon_o V_s$$
$$0 + 0 = 4\pi\varepsilon_o V_s$$

Apparently works at all only if $V_s = 0$.

$$r \rightarrow 0 \qquad \frac{q_o}{\sqrt{z_o'^2}} + \frac{q_i}{\sqrt{z_i'^2}} = 0$$
$$\frac{q_i}{|z_i'|} = -\frac{q_o}{|z_o'|}$$

Example: a charge q_o suspended distance z_o above a flat conducting surface that's held at voltage V_s . Find expression for V anywhere above the conducting surface.



Example: a charge q_o suspended distance z_o above a flat conducting surface that's held at voltage V_s . Find expression for V anywhere above the conducting surface.

Digression: Force between *q_o* and Surface

Image charge distance and magnitude

$$z_i' = -z_o' \qquad q_i = -q_o$$

Field everywhere above the plane is as if there were these two charges

$$\vec{F}_{i\to o} = \frac{1}{4\pi\varepsilon_o} \frac{q_i q_o}{\left(z_o - z_i\right)^2} \,\hat{z}$$

$$\vec{F}_{i \to o} = -\frac{1}{4\pi\varepsilon_o} \frac{q_o^2}{(2z_o)^2} \hat{z}$$

 $\vec{F}_{o \to surface} = \frac{1}{4\pi\varepsilon_o} \frac{q_o^2}{(2z_o)^2} \hat{z}$



Example: a charge q_o suspended distance z_o above a flat conducting surface that's held at voltage V_s . Find expression for V anywhere above the conducting surface.

Digression: Work of moving *q_o* into place (and arranging surface charge)



$$z_i' = -z_o' \qquad q_i = -q_o$$

Field everywhere above the plane is as if there were these two charges

$$\vec{F}_{surface \to o} = -\frac{1}{4\pi\varepsilon_o} \frac{q_o^2}{(2z_o)^2} \hat{z}$$

$$W = \int_{\infty}^{z_o} \vec{F}_{you \to o} \cdot d\vec{l} = -\int_{\infty}^{z_o} \vec{F}_{surface \to o} \cdot d\vec{l} = \int_{z_o}^{\infty} \vec{F}_{surface \to o} \cdot d\vec{l}$$

$$W = \int_{z_o}^{\infty} -\frac{1}{4\pi\varepsilon_o} \frac{q_o^2}{(2z)^2} \,\hat{z} \cdot d\vec{l} = \int_{z_o}^{\infty} -\frac{1}{4\pi\varepsilon_o} \frac{q_o^2}{(2z)^2} \,dz$$

$$W = -\frac{q_o^2}{16\pi\varepsilon_o} \int_{z_o}^{\infty} \frac{1}{z^2} dz = \frac{q_o^2}{16\pi\varepsilon_o} \frac{1}{z} \Big|_{z_o}^{\infty} = -\frac{q_o^2}{16\pi\varepsilon_o} \frac{1}{z_o}$$



Example: a charge q_o suspended distance z_o above a flat conducting surface that's held at voltage V_s . Find expression for V anywhere above the conducting surface.



Example: a charge q_o suspended distance z_o above a flat conducting surface that's held at voltage V_s . Find expression for V anywhere above the conducting surface.



Surface Charge Density

$$\frac{q_o}{4\pi\varepsilon_o} \left(\frac{1}{\sqrt{x^2 + y^2 + (z - z'_o)^2}} - \frac{1}{\sqrt{x^2 + y^2 + (z + z'_o)^2}} \right) + V_s = V(\vec{r})$$

From a Gaussian pill-box: $E_n = \frac{\sigma}{m}$ n denotes component perpendicular to surface In terms of V: $-\frac{\partial V}{\partial n} = E_n = \frac{\sigma}{\varepsilon}$ So, $\sigma = -\varepsilon_o \frac{\partial V}{\partial n}$ In this case, $\sigma = -\varepsilon_o \frac{\partial V}{\partial z} \bigg|_{z=0} - \frac{q_o}{4\pi} \frac{\partial}{\partial z} \bigg(\frac{1}{\sqrt{x^2 + y^2 + (z - z_o')^2}} - \frac{1}{\sqrt{x^2 + y^2 + (z + z_o')^2}} \bigg) \bigg|_{z=0}$ â $z_{q}' \overset{\mathsf{Y}_{0}}{\vec{r}} \overset{\mathbf{v}_{1}}{\vec{r}} = -\frac{q_{o}}{4\pi} \left(\frac{-1}{2} \frac{2(z-z_{o}')}{(x^{2}+y^{2}+(z-z_{o}')^{2})^{\frac{3}{2}}} - \frac{-1}{2} \frac{2(z+z_{o}')}{(x^{2}+y^{2}+(z+z_{o}')^{2})^{\frac{3}{2}}} \right)$ $\hat{y} = -\frac{q_o}{4\pi} \left(\frac{(z'_o)}{(x^2 + v^2 + (z')^2)^{\frac{3}{2}}} + \frac{(z'_o)}{(x^2 + v^2 + (z')^2)^{\frac{3}{2}}} \right)$ $\sigma = -\frac{q_o}{2\pi} \frac{z_o}{\left(x^2 + v^2 + {z'_o}^2\right)^{\frac{3}{2}}}$

Surface Charge

$$q_{surf} = \int \sigma da$$

$$q_{surf} = \int -\frac{q_o}{4\pi} \frac{2z'_o}{\left(x^2 + y^2 + {z'_o}^2\right)^{3/2}} dx dy = -\frac{q_o}{4\pi} \int \frac{2z'_o}{\left(s^2 + {z'_o}^2\right)^{3/2}} s d\phi ds = -q_o \int \frac{z'_o}{\left(s^2 + {z'_o}^2\right)^{3/2}} \frac{1}{2} ds^2$$

$$q_{surf} = -q_o \left(-\frac{z'_o}{\left(s^2 + {z'_o}^2\right)^{1/2}} \right|_{s=0}^{s=\infty} \right) = q_o \left(0 - \frac{z'_o}{\left(z'_o^2\right)^{1/2}} \right) = -q_o$$

$$z_q \stackrel{q_o}{i} \stackrel{q_o}{r} \stackrel{q}{r} \stackrel{q}{r}$$

General Approach

•Draw picture

•Appeal to symmetry (and intuition about mirrors)

•Apply the condition $\sum \frac{q_o}{n_e} + \sum \frac{q_i}{n_e} = const$ on conductor •See what you've got to do to remove dependence on the observation location on conductor.

1)Mathematically, you've got 3 free parameters: the constant, the image charge's value, and the image charge's location.

2) Since the relation should be true for *all* observation locations on the conducting surface, choose easy ones to help you determine the three parameter's values.

•If you've got a solution that works for the boundary and satisfies Poisson's equation, then you've got *the* solution.



Exercise: where and what are the image charges?

What is the force on charge -2q_o?

3d

d

What is the work of bringing it into position, once $-2q_o$ $q_o's$ already there?

Exercise: charge and bent conductor, what's the equivalent charge distribution?



V= 0

What's the potential expression in front of the conductor? What's the charge density on each face of the conductor?



Appeal to symmetry (and intuition about mirrors)



•Appeal to symmetry (and intuition about mirrors)

$$\frac{\hat{q}_{o}}{\hat{v}_{o}} + \frac{\hat{q}_{i}}{\hat{v}_{i}} = 4\pi\varepsilon_{o}V(\vec{r})$$

$$\frac{\hat{q}_{o}}{\hat{v}_{o}} + \frac{\hat{q}_{i}}{\hat{v}_{i}} = 4\pi\varepsilon_{o}V(\vec{r})$$

$$\frac{\hat{q}_{o}}{\sqrt{(\vec{r} - \vec{r}_{o}')^{2}}} + \frac{\hat{q}_{i}}{\sqrt{(\vec{r} - \vec{r}_{i}')^{2}}} = 4\pi\varepsilon_{o}V(\vec{r})$$

$$\frac{\hat{q}_{o}}{\sqrt{r^{2} + r_{o}'^{2} - 2rr_{o}'\cos\theta_{o}}} + \frac{\hat{q}_{i}}{\sqrt{r^{2} + r_{i}'^{2} - 2rr_{i}'\cos\theta_{i}}} = 4\pi\varepsilon_{o}V(\vec{r})$$

$$V(R) = 0$$

$$\frac{\hat{v}_{i}}{\hat{v}_{i}} + \frac{\hat{v}_{o}}{\sqrt{r^{2} + a^{2} - 2ra\cos\theta_{o}}} + \frac{\hat{q}_{i}}{\sqrt{r^{2} + b^{2} - 2rb\cos\theta_{i}}} = 4\pi\varepsilon_{o}V(\vec{r})$$

$$\frac{\hat{v}_{i}}{\hat{v}_{i}} + \frac{\hat{v}_{o}}{\sqrt{r^{2} + a^{2} - 2ra\cos\theta_{o}}} + \frac{\hat{v}_{o}}{\sqrt{r^{2} + b^{2} - 2rb\cos\theta_{i}}} = 4\pi\varepsilon_{o}V(\vec{r})$$

ŷ



$$\frac{q_o}{\sqrt{r^2 + a^2 - 2ra\cos\theta_o}} + \frac{q_i}{\sqrt{r^2 + b^2 - 2rb\cos\theta_i}} = 4\pi\varepsilon_o V(\vec{r})$$

by the condition
$$\sum \frac{q_o}{\eta_o} + \sum \frac{q_i}{\eta_i} = const \text{ on conductor}$$
$$\frac{q_o}{\sqrt{R^2 + a^2 - 2Ra\cos\theta_o}} + \frac{q_i}{\sqrt{R^2 + b^2 - 2Rb\cos\theta_i}} = 0$$

•See what you've got to do to remove dependence on the observation location on conductor.

3 free parameters: the constant, the image charge's value, and the image charge's location.

The relation should be true for *all* observation locations on the conducting surface, so choose easy ones to help you determine the three parameter's values.



$$\frac{q_o}{\sqrt{r^2 + a^2 - 2ra\cos\theta_o}} + \frac{q_i}{\sqrt{r^2 + b^2 - 2rb\cos\theta_i}} = 4\pi\varepsilon_o V(\vec{r})$$

by the condition
$$\sum \frac{q_o}{\eta_o} + \sum \frac{q_i}{\eta_i} = const \text{ on conductor}$$
$$\frac{q_o}{\sqrt{R^2 + a^2 - 2Ra\cos\theta_o}} + \frac{q_i}{\sqrt{R^2 + b^2 - 2Rb\cos\theta_i}} = 0$$

•See what you've got to do to remove dependence on the observation location on conductor.

3 free parameters: the constant, the image charge's value, and the image charge's location.

The relation should be true for *all* observation locations on the conducting surface, so choose easy ones to help you determine the three parameter's values.

 \hat{y} Top of sphere:

$$\frac{q_o}{(a-R)} + \frac{q_i}{(R-b)} = 0$$



$\frac{q_o}{\sqrt{2}} + \frac{1}{\sqrt{2}}$	$\frac{q_i}{1^2 - 2 - 1} = 4\pi \varepsilon_o V(\vec{r})$
$\sqrt{r^2 + a^2 - 2ra\cos\theta_o} \sqrt{r^2 + a^2}$	$b^2 - 2rb\cos\theta_i$
ply the condition $\sum \frac{q_o}{d} + \sum q_o$	$\frac{q_i}{dt} = const$ on conductor
- <i>n</i> _o -	r.
$\underline{q_o}$	$+ \frac{q_i}{q_i} = 0^{4}$
$\sqrt{R^2 + a^2 - 2Ra\cos\theta_o}$	$\sqrt{R^2 + b^2 - 2Rb\cos\theta_i}$

•See what you've got to do to remove dependence on the observation location on conductor.

3 free parameters: the constant, the image charge's value, and the image charge's location.

The relation should be true for *all* observation locations on the conducting surface, so choose easy ones to help you determine the three parameter's values.

 \hat{y} Top of sphere:

$$\frac{q_o}{(a-R)} + \frac{q_i}{(R-b)} = 0$$

Bottom of sphere:

$$\frac{q_o}{(a+R)} + \frac{q_i}{(R+b)} = 0$$

•See what you've got to do to remove dependence on the observation location on conductor.

3 free parameters: the constant, the image charge's value, and the image charge's location.

The relation should be true for *all* observation locations on the conducting surface, so choose easy ones to help you determine the three parameter's values.

Top of sphere:

 θ_{o}

Ŕ

R

Ħ,

q_o

 \vec{r}_{o}

 $z_{q}' = a$

 $z_i'=b^{q_i}$

 \vec{r}

 \vec{r}

V(R)

 \widehat{x}

Bottom of sphere:

$$\frac{q_o}{(a-R)} + \frac{q_i}{(R-b)} = 0 \qquad \qquad \frac{q_o}{(a+R)} + \frac{q_i}{(R+b)} = 0$$

Two equations, two unknowns (qi, b)

$$-q_{o} \frac{(R-b)}{(a-R)} = q_{i} = -q_{o} \frac{(R+b)}{(a+R)}$$
$$(R-b)(a+R) = (R+b)(a-R)$$
$$(R-ab-Rb+R^{2}) = (Ra-R^{2}+ab-Rb)$$
$$(Ra-ab-Rb+R^{2}) = (Ra-R^{2}+ab-Rb)$$
$$(-ab+R^{2}) = (-R^{2}+ab)$$
$$b = \frac{R^{2}}{a}$$

•See what you've got to do to remove dependence on the observation location on conductor.

3 free parameters: the constant, the image charge's value, and the image charge's location.

The relation should be true for *all* observation locations on the conducting surface, so choose easy ones to help you determine the three parameter's values.

$$b = \frac{R^2}{a} \qquad -q_o \frac{(R-b)}{(a-R)} = q_i = -q_o \frac{(R+b)}{(a+R)}$$
$$q_i = -q_o \frac{(R+R^2/a)}{(a+R)}$$
$$q_i = -q_o \frac{R}{a} \frac{(1+R)}{(1+R)}$$
$$q_i = -q_o \frac{R}{a}$$





•If you've got a solution that works for the boundary and satisfies Poisson's equation, then you've got *the* solution.



 \hat{x}



ŷ

•Apply the condition $\sum \frac{q_o}{\mathbf{v}_o} + \sum \frac{q_i}{\mathbf{v}_i} = const \text{ on conductor}$ $\frac{q_o}{\sqrt{R^2 + a^2 - 2Ra\cos\theta_o}} + \frac{q_i}{\sqrt{R^2 + b^2 - 2Rb\cos\theta_i}} = 4\pi\varepsilon_o V(R) = 0$

•See what you've got to do to remove dependence on the observation location on conductor.

3 free parameters: the constant, the image charge's value, and the image charge's location.

The relation should be true for *all* observation locations on the conducting surface, so choose easy ones to help you determine the three parameter's values.



Exercise: charge and flat/curved conductor, what's the equivalent q charge distribution?



Fri.	3.2 Images T4 Relaxation Method	
Mon. Wed. Thurs	3.4.14.2 Multipole Expansion 3.4.34.4 Multipole Expansion	HW4